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MULTIGRID FOR THE MORTAR FINITE ELEMENT METHOD*

JAYADEEP GOPALAKRISHNAN[†] AND JOSEPH E. PASCIAK[†]

Abstract. A multigrid technique for uniformly preconditioning linear systems arising from a mortar finite element discretization of second order elliptic boundary value problems is described and analyzed. These problems are posed on domains partitioned into subdomains, each of which is independently triangulated in a multilevel fashion. The multilevel mortar finite element spaces based on such triangulations (which need not align across subdomain interfaces) are in general not nested. Suitable grid transfer operators and smoothers are developed which lead to a variable V-cycle preconditioner resulting in a uniformly preconditioned algebraic system. Computational results illustrating the theory are also presented.

Key words. mortar, finite element method, multigrid, V-cycle, preconditioning, domain decomposition

AMS subject classifications. 65F10, 65N55, 65N30

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1. Introduction. The mortar finite element method is a nonconforming domain decomposition technique tailored to handle problems posed on domains that are partitioned into independently triangulated subdomains. The meshes on different subdomains need not align across subdomain interfaces. The flexibility this technique offers by allowing substructures of a complicated domain to be meshed independently of each other is well recognized. In this paper we consider preconditioned iteration for the solution of the resulting algebraic system. Our preconditioner is a nonvariational multigrid procedure.

The mortar finite element discretization is a discontinuous Galerkin approximation. The functions in the approximation subspaces have jumps across subdomain interfaces and are standard finite element functions when restricted to the subdomains. The jumps across subdomain interfaces are constrained by conditions associated with one of the two neighboring meshes. Bernardi, Maday, and Patera (see [3, 4]) proved the coercivity of the associated bilinear form on the mortar finite element space, thus implying existence and uniqueness of solutions to the discrete problem. They also showed that the mortar finite element method is as accurate as the usual finite element method. Recently, stability and convergence estimates for an hp version of the mortar finite element method were proved [19].

When each subdomain has a multilevel mesh, preconditioners for the linear system arising from the mortar discretization can be developed by multilevel techniques. A hierarchical preconditioner with conditioning which grows like the square of the number of levels is described in [11]. Substructuring preconditioners have also been developed for the mortar method [1]. In this paper, we show that a variable V-cycle may be used to develop a preconditioned system whose condition number remains bounded independently of the number of levels. A multigrid preconditioner for the

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mortar method viewed as a saddle point problem is described in [5]. In contrast, our algorithm gives rise to a positive definite preconditioned system.

One of the difficulties in constructing a multigrid preconditioner for the mortar finite element method arises due to the fact that the multilevel mortar finite element spaces are, in general, not nested. Multigrid theory for nonnested spaces [7] may be employed to construct a variable V-cycle preconditioner, provided a suitable prolongation operator can be designed. We construct such a prolongation operator and prove that it satisfies the “regularity and approximation” property (Condition (C.2)) required for application of the multigrid theory.

The next difficulty is in the design of a smoother. Our smoother is based on the point Jacobi method. Its analysis is nonstandard since the constraints at subdomain interface gives rise to mortar basis functions with nonlocal support. We prove that these basis functions decay exponentially away from their nodal vertex. This leads to a strengthened Cauchy–Schwarz inequality which is used to verify the smoothing hypothesis (Condition (C.1)).

The remainder of the paper is organized as follows. Section 2 introduces most of the notation in the paper. Section 3 describes the multilevel mortar finite element spaces. In section 4 the variable V-cycle multigrid algorithm is given and the main result (Theorem 4.1) is stated and proved. Section 5 provides proofs of some technical lemmas. Implementation issues are considered in section 6 while the results of numerical experiments illustrating the theory are given in section 7.

2. Preliminaries. In this section, we provide some preliminaries and notation which will be used in the remainder of the paper. In addition, we describe the continuous problem and impose an assumption on the regularity of its solution.

Let Ω be an open subset of the plane. For nonnegative integers s , the Sobolev space $H^s(\Omega)$ (see [10, 14]) is the set of functions in $L^2(\Omega)$ with distributional derivatives up to order s also in $L^2(\Omega)$. If s is a positive real number between nonnegative integers m and $m + 1$, $H^s(\Omega)$ is the space obtained by interpolation (by the real method [16]) between $H^m(\Omega)$ and $H^{m+1}(\Omega)$. The Sobolev norm on $H^s(\Omega)$ is denoted by $\|\cdot\|_{s,\Omega}$ and the corresponding Sobolev seminorm is denoted by $|\cdot|_{s,\Omega}$. For $\phi \in H^s(\Omega)$ and a segment γ contained in $\bar{\Omega}$, the trace of ϕ on γ is denoted by $\phi|_\gamma$. We will often write $\|\phi\|_{r,\gamma}$ and $|\phi|_{r,\gamma}$ for the $H^r(\gamma)$ norm and seminorm, respectively, of the trace $\phi|_\gamma$.

Assume that Ω is connected and that its boundary, $\partial\Omega$, is polygonal. Let $\partial\Omega$ be split into $\partial\Omega_D$ and $\partial\Omega_N$ such that $\partial\Omega = \bar{\partial\Omega}_N \cup \bar{\partial\Omega}_D$ and $\partial\Omega_N \cap \partial\Omega_D$ is empty and assume that $\partial\Omega_D$ has nonzero measure. Denote by \mathcal{V} the subspace of the Sobolev space $H^1(\Omega)$ consisting of functions in $H^1(\Omega)$ whose trace on $\partial\Omega_D$ is zero. Denote by \mathcal{V}' the dual of the normed linear space \mathcal{V} . The dual norm $\|\cdot\|_{-1,\Omega}$ is defined by

$$\|u\|_{-1,\Omega} = \sup_{\phi \in \mathcal{V}} \frac{\langle u, \phi \rangle}{\|\phi\|_{1,\Omega}},$$

where $\langle \cdot, \cdot \rangle$ denotes the duality pairing. Note that $L^2(\Omega)$ is contained in \mathcal{V}' if we identify the functional $\langle v, \phi \rangle = (v, \phi)$ for all $v \in L^2(\Omega)$. Here (\cdot, \cdot) denotes the inner-product in $L^2(\Omega)$. For $-1 < s < 0$, $\|\cdot\|_{s,\Omega}$ is the norm on the space defined by interpolation between \mathcal{V}' and $L^2(\Omega)$.

We seek an approximate solution to the problem

$$(2.1) \quad A(U, \phi) = F(\phi) \text{ for all } \phi \in \mathcal{V},$$

where $A(\cdot, \cdot)$ is bilinear form on $\mathcal{V} \times \mathcal{V}$ defined by

$$A(u, v) = \int_{\Omega} \nabla u \cdot \nabla v \, dx,$$

and F is a given continuous linear functional on $H^1(\Omega)$. This problem has a unique solution. For the mortar finite element method, we restrict our attention to F of the form

$$(2.2) \quad F(v) = \int_{\Omega} f v \, dx$$

for $f \in L^2(\Omega)$. This is the variational form of the boundary value problem

$$(2.3) \quad -\Delta U = f \text{ in } \Omega,$$

$$(2.4) \quad U = 0 \text{ on } \partial\Omega_D,$$

$$(2.5) \quad \frac{\partial U}{\partial n} = 0 \text{ on } \partial\Omega_N.$$

Although our results are stated for this model problem, extension to more general second order elliptic partial differential equations with more general boundary conditions are straightforward.

We will need to assume some regularity for solutions of Problem (2.1). We formalize it here into Assumption (A.1).

(A.1) There exists a β in the interval $(1/2, 1]$ for which

$$\|U\|_{1+\beta, \Omega} \leq C \|F\|_{-1+\beta, \Omega}$$

holds for solutions U to the problem (2.1).

This is known to hold for a wide class of domains [15]. Note that we do not require full elliptic regularity ($\beta = 1$ case).

3. The mortar finite element method. In this section, we first provide notation for subdomains and triangulations. Next multilevel mortar finite element spaces are introduced and the mortar finite element problem is defined.

Partition Ω into nonoverlapping polygonal subdomains Ω_i , $i = 1, \dots, K$. The interface $\Gamma = \cup_{i=1}^K \partial\Omega_i \setminus \partial\Omega$ is broken into a set of disjoint open straight line segments γ_k , each of which is contained in $\partial\Omega_i \cap \partial\Omega_j$ for some i and j . The collection of these edges will be denoted by Z , i.e., $Z = \{\gamma_1, \gamma_2, \dots, \gamma_L\}$.

Each Ω_i is triangulated to produce a quasi-uniform mesh \mathcal{T}_1^i of size h_1 . The triangulations generally do not align at the subdomain interfaces. We assume that the endpoints of each interface segment in Z are vertices of \mathcal{T}_1^p and \mathcal{T}_1^q where p and q are such that $\gamma \subset \partial\Omega_p \cap \partial\Omega_q$. Denote the global mesh $\cup_i \mathcal{T}_1^i$ by \mathcal{T}_1 . To set up the multigrid algorithm, we need a sequence of refinements of \mathcal{T}_1 . We refine the triangulation \mathcal{T}_1 to produce \mathcal{T}_2 by splitting each triangle of \mathcal{T}_1 into four triangles by joining the midpoints of the edges of the triangle. The triangulation \mathcal{T}_2 is then quasi-uniform of size $h_2 = h_1/2$. Repeating this process, we get a sequence of triangulations \mathcal{T}_k , $k = 1, \dots, J$, each quasi-uniform of size $h_k = h_1/2^{k-1}$.

We next define the mortar finite element spaces following [2, 3, 4, 19] (our notation is close to that in [19]). First, we define spaces $\tilde{\mathcal{V}}$ and \tilde{M}_k by

$$(3.1) \quad \tilde{\mathcal{V}} = \{v : v|_{\Omega_i} \in H^1(\Omega_i) \quad \text{for all } i = 1, \dots, K, v = 0 \text{ on } \partial\Omega_D\}$$

and

$$\widetilde{M}_k = \{v \in \widetilde{\mathcal{V}} : v \text{ is linear on each triangle of } \mathcal{T}_k\}.$$

Throughout this paper we will use piecewise linear finite element spaces for convenience of notation. The results extend to higher order finite elements without difficulty [13].

For every straight line segment $\gamma \in Z$, there is an i and j such that $\gamma \subseteq \partial\Omega_i \cap \partial\Omega_j$. Assign one of i and j to be the *mortar* index, $M(\gamma)$, and the other then is the *nonmortar* index, $NM(\gamma)$. Let $\Omega_{M(\gamma)}$ denote the mortar domain of γ and $\Omega_{NM(\gamma)}$ be the nonmortar domain of γ . For every $u \in \widetilde{\mathcal{V}}$ define u_γ^M and u_γ^{NM} to be the trace of $u|_{\Omega_{M(\gamma)}}$ on γ and the trace of $u|_{\Omega_{NM(\gamma)}}$ on γ , respectively.

We now define two discrete spaces $S_k(\gamma)$ and $W_k(\gamma)$ on an interface segment γ . Every $\gamma \in Z$ can be divided into subintervals in two ways: by the vertices of the mesh in the mortar domain of γ and by those of the nonmortar domain of γ . Consider $\bar{\gamma}$ as partitioned into subintervals by the vertices of the triangulation on nonmortar side. Let these vertices be denoted by $x_{k,\gamma}^i$, $i = 0, \dots, N$. Denote the subintervals $[x_{k,\gamma}^{i-1}, x_{k,\gamma}^i]$ by $\omega_{k,i}$, $i = 1, \dots, N$, where $\omega_{k,1}$ and $\omega_{k,N}$ are the subintervals that are at the ends of γ . The discrete space $S_k(\gamma)$ is defined as follows:

$$S_k(\gamma) = \left\{ \begin{array}{l} v \text{ is linear on each } \omega_{k,i}, \ i = 1, \dots, N, \\ v : \ v \text{ is constant on } \omega_{k,1} \text{ and on } \omega_{k,N}, \\ \text{and } v \text{ is continuous on } \gamma. \end{array} \right\}.$$

We also define the space $W_k(\gamma)$ by

$$W_k(\gamma) = \left\{ \begin{array}{l} v \text{ is linear on each } \omega_{k,i}, \ i = 1, \dots, N, \\ v : \ v \text{ vanishes at endpoints of } \gamma, \text{ namely } x_{k,\gamma}^0 \text{ and } x_{k,\gamma}^N, \\ \text{and } v \text{ is continuous on } \gamma. \end{array} \right\}.$$

The multilevel mortar finite element spaces M_k , $k = 1, \dots, J$, are now defined by

$$(3.2) \quad M_k = \left\{ u \in \widetilde{M}_k : \begin{array}{l} \text{on each } \gamma \in Z, \int_\gamma (u_\gamma^M - u_\gamma^{NM}) \chi \, ds = 0 \\ \text{for all } \chi \in S_k(\gamma). \end{array} \right\}.$$

The ‘‘mortaring’’ is done by constraining the jump across interfaces by the integral equality above. We will call this constraint the *weak continuity* of functions in M_k . Note that though the spaces $\{\widetilde{M}_k\}$ are nested,

$$\widetilde{M}_1 \subseteq \dots \subseteq \widetilde{M}_k \subseteq \widetilde{M}_{k+1} \subseteq \dots \subseteq \widetilde{M}_J,$$

the multilevel spaces $\{M_k\}$ are generally nonnested.

We next state the error estimates for the mortar finite element method. The mortar finite element approximation of the solution U of Problem (2.1) (with F given by (2.2)) is the function $U_k \in M_k$ satisfying

$$(3.3) \quad \widetilde{A}(U_k, \phi) = \int_\Omega f \phi \, dx \quad \text{for all } \phi \in M_k,$$

where $\widetilde{A}(u, v)$ is the bilinear form on $\widetilde{\mathcal{V}} \times \widetilde{\mathcal{V}}$ defined by

$$\widetilde{A}(u, v) = \sum_{i=1}^K \int_{\Omega_i} \nabla u \cdot \nabla v \, dx.$$

It is shown in [3] that

$$\|u\| \leq C\tilde{A}(u, u) \quad \text{for all } u \in \tilde{M}_k,$$

where $\|v\|^2 = \sum_{i=1}^K \|v\|_{1,\Omega_i}^2$. Here and in the remainder of this paper, we will use C to denote a generic constant independent of h_k which can be different at different occurrences. It follows that (3.3) has a unique solution. It is also known (see [3]) that the mortar finite element approximation satisfies

$$(3.4) \quad \|u - U_k\| \leq Ch_k^\beta \|u\|_{1+\beta,\Omega}.$$

We now define a projection, $\Pi_{k,\gamma} : L^2(\gamma) \rightarrow W_k(\gamma)$, which will be very useful in our analysis. For $u \in L^2(\gamma)$, it can be shown [4] that there exists a unique $v \in W_k(\gamma)$ satisfying

$$(3.5) \quad \int_\gamma v\chi \, ds = \int_\gamma u\chi \, ds \quad \text{for all } \chi \in S_k(\gamma).$$

We define $\Pi_{k,\gamma}u$ to be v . This projection is known to be stable in $L^2(\gamma)$ and $H_0^1(\gamma)$, i.e.,

$$(3.6) \quad \|\Pi_{k,\gamma}u\|_{0,\gamma} \leq C \|u\|_{0,\gamma} \quad \text{for all } u \in L^2(\gamma) \text{ and}$$

$$(3.7) \quad \|\Pi_{k,\gamma}u\|_{1,\gamma} \leq C \|u\|_{1,\gamma} \quad \text{for all } u \in H_0^1(\gamma),$$

under some weak assumptions on meshes (see [19]) which hold for the meshes defined above.

The projector $\Pi_{k,\gamma}$ is clearly related to the weak continuity condition. Let $\{y_k^j\}$ denote the nodes of \mathcal{T}_k and the operator $\mathcal{E}_{k,\gamma} : \tilde{\mathcal{V}} \rightarrow \tilde{M}_k$ be defined by (see also Figure 3.1(a)–(d))

$$(3.8) \quad \mathcal{E}_{k,\gamma}\tilde{u}(y_k^j) = \begin{cases} (\Pi_{k,\gamma}(\tilde{u}_\gamma^M - \tilde{u}_\gamma^{NM})) (y_k^j) & \text{if } y_k^j \in \gamma \cap \bar{\Omega}_{NM}(\gamma), \\ 0 & \text{otherwise.} \end{cases}$$

It is easy to see that if \tilde{u} is in \tilde{M}_k then $u = \tilde{u} + \sum_{\gamma \in Z} \mathcal{E}_{k,\gamma}\tilde{u}$ is an element of M_k .

We next define a basis for M_k . Let $\{\tilde{\phi}_k^i : i = 1, \dots, \tilde{N}_k\}$ be the nodal basis for \tilde{M}_k . There is more than one basis element associated with a node which appears in multiple subdomains. The basis for M_k consists of functions of the form

$$(3.9) \quad \phi_k^i = \tilde{\phi}_k^i + \sum_{\gamma \in Z} \mathcal{E}_{k,\gamma}(\tilde{\phi}_k^i).$$

For every vertex y_k^l located in the open segment $\gamma \in Z$ and belonging to the nonmortar side mesh, the corresponding ϕ_k^l as defined above is zero. Every remaining vertex y_k^l leads to a nonzero ϕ_k^l since ϕ_k^l and $\tilde{\phi}_k^l$ have the same nonzero value at y_k^l . Also, the values of ϕ_k^l and $\tilde{\phi}_k^l$ at all nodes which are not nodes from nonmortar mesh lying in the interior of some $\gamma \in Z$ are the same. This implies that nonzero functions in $\{\phi_k^i\}$ are linearly independent. It is not difficult to check that these also form a basis for M_k . Since at y_k^l , ϕ_k^l is one and all other ϕ_k^i for $i \neq l$ are zero, these functions, in fact, form a nodal basis. Denote by N_k the total number of nonzero ϕ_k^i . We now reindex

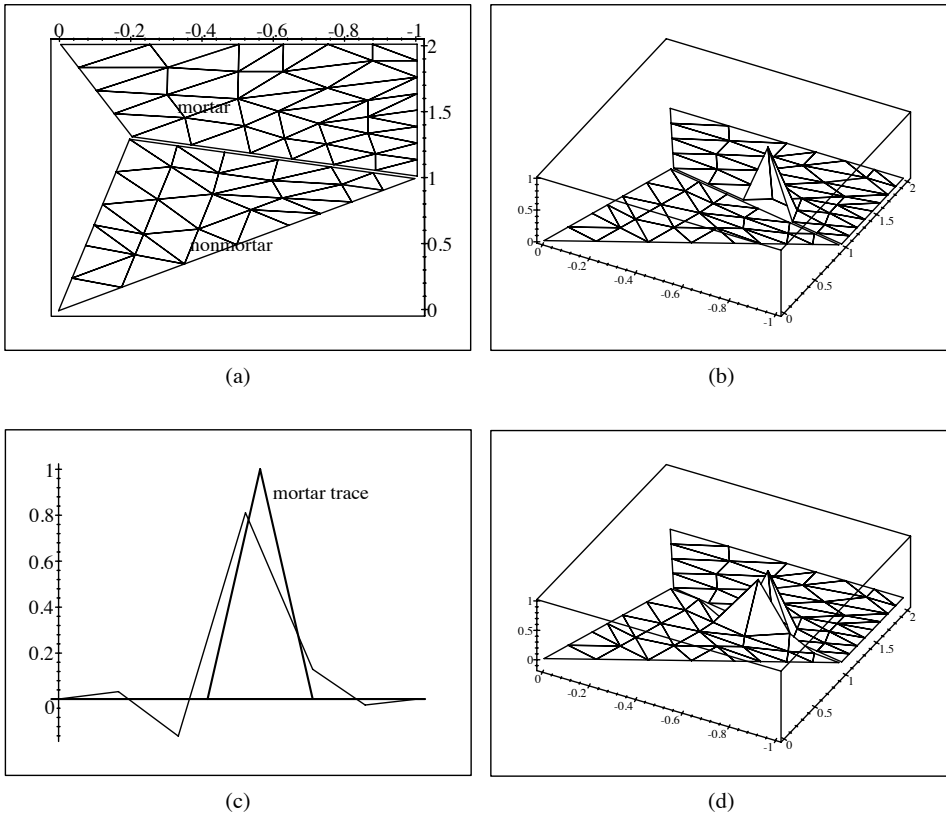


FIG. 3.1. Illustrating the action of $\mathcal{E}_{k,\gamma}$: (a) Two subdomains with meshes that do not align at interface. (b) A discontinuous \tilde{u} , which is 1 on a mortar node and 0 on the remaining nodes. (c) The thick line shows \tilde{u}_γ^M and the thin line shows $\Pi_{k,\gamma}(\tilde{u}_\gamma^M - 0)$. (d) Plot shows $\tilde{u} + \mathcal{E}_{k,\gamma}\tilde{u}$. $\mathcal{E}_{k,\gamma}\tilde{u}$ is formed by extending $\Pi_{k,\gamma}\tilde{u}_\gamma^M$ as described by (3.8).

$\{\tilde{\phi}_k^i : i = 1, \dots, \tilde{N}_k\}$ in such a way that every nonzero ϕ_k^i is in $\{\phi_k^i : i = 1, \dots, N_k\}$. We also reindex $\{y_k^i\}$ in this new ordering.

Now that we have a nodal basis for M_k , we may speak of the corresponding vertices of \mathcal{T}_k as degrees of freedom for M_k . Consider an interface segment $\gamma \in Z$. All vertices on $\bar{\gamma}$ are degrees of freedom except (i) those on $\partial\Omega_D$ and (ii) those on γ and from the nonmortar mesh. These are the vertices $y_k^i, i = 1, \dots, N_k$.

4. Multigrid algorithm for the mortar FEM. We will apply multigrid theory for nonnested spaces [7] to construct a variable V-cycle preconditioner. Before giving the algorithm, we define a prolongation operator and smoother. Later in this section, we will prove that our algorithm gives a preconditioner which results in a preconditioned system with uniformly bounded condition number.

First, let us establish some notation: A_k will denote the operator on M_k , generated by the form $\tilde{A}(\cdot, \cdot)$, i.e., A_k is defined by

$$(A_k u, v) = \tilde{A}(u, v) \quad \text{for all } u, v \in M_k.$$

The largest eigenvalue of A_k is denoted by λ_k . For each basis element ϕ_k^i , we define

$M_k^i, i = 1, \dots, N_k$, to be the one dimensional subspace of M_k spanned by ϕ_k^i . Then

$$M_k = \sum_{i=1}^{N_k} M_k^i$$

provides a direct sum decomposition of M_k .

4.1. Smoothing and prolongation operators. We will use a smoother R_k given by a scaled Jacobi method, i.e.,

$$(4.1) \quad R_k = \alpha \sum_{i=1}^{N_k} A_{k,i}^{-1} Q_k^i,$$

where α is a positive constant to be chosen later. Here, $A_{k,i} : M_k^i \rightarrow M_k^i$ and $Q_k^i : L^2(\Omega) \rightarrow M_k^i$ are defined by

$$(A_{k,i}v, \chi) = A(v, \chi) \quad \text{for all } \chi \in M_k^i,$$

and

$$(Q_k^i v, \chi) = (v, \chi) \quad \text{for all } \chi \in M_k^i,$$

respectively. R_k is symmetric in the (\cdot, \cdot) inner-product. The following property of R_k will be proved in section 5.

(C.1) There exists a positive number C_R independent of k such that

$$(4.2) \quad \frac{\|u\|_{0,\Omega}^2}{\lambda_k} \leq C_R (R_k u, u) \quad \text{for all } u \in M_k.$$

In addition, $I - R_k A_k$ is nonnegative.

We now define ‘‘prolongation operators’’ $I_k : M_{k-1} \rightarrow M_k$, for $k = 2, \dots, J$. Clearly, $I_k u$ needs to satisfy the weak continuity constraint (see Definition 3.2). We define $I_k u$ by

$$(4.3) \quad I_k u = u + \sum_{\gamma \in Z} \mathcal{E}_{k,\gamma}(u).$$

In the next section we show that I_k satisfies the following inequality.

(C.2) There exists a constant C_β independent of k such that

$$|A_k((I - I_k P_{k-1})u, u)| \leq C_\beta^2 \left(\frac{\|A_k u\|_{0,\Omega}^2}{\lambda_k} \right)^{\beta/2} \left(\tilde{A}(u, u) \right)^{1-\beta/2}$$

for all u in M_k .

Here P_k is the \tilde{A} -adjoint of I_k , i.e., $P_k : M_{k+1} \rightarrow M_k, k = 1, \dots, J - 1$, satisfies

$$\tilde{A}(P_k u, \phi) = \tilde{A}(u, I_{k+1} \phi) \quad \text{for all } \phi \in M_k.$$

Condition (C.2) is verified using the regularity of the underlying partial differential equation.

4.2. The algorithm. Let $m(k)$, $k = 1, \dots, J$, be positive integers depending on k and $P_{k-1}^0 : M_k \rightarrow M_{k-1}$ be defined by

$$(P_{k-1}^0 u, v) = (u, I_k v) \quad \text{for all } u \in M_k \quad \text{and} \quad v \in M_{k-1}.$$

The variable V-cycle preconditioner B_k for $k = 1, \dots, J$ is defined as follows.

ALGORITHM 4.2.

1. For $k = 1$, set $B_1 = A_1^{-1}$.
2. For $k = 2, \dots, J$, $B_k g$ for $g \in M_k$ is defined recursively by the following five steps.

- (a) Set $x^0 = 0$.
- (b) Define x^l , for $l = 1, \dots, m(k)$ by

$$x^l = x^{l-1} + R_k(g - A_k x^{l-1}).$$

- (c) Set $y^0 = x^{m(k)} + I_k q$, where q is given by

$$q = B_{k-1} P_{k-1}^0 (g - A_k x^{m(k)}).$$

- (d) Define y^l for $l = 1, \dots, m(k)$ by

$$y^l = y^{l-1} + R_k(g - A_k y^{l-1}).$$

- (e) Set $B_k g \equiv y^{m(k)}$.

We make the usual assumption on $m(k)$ (cf. [7]).

(A.2) The number of smoothings $m(k)$ increases as k decreases in such a way that

$$\beta_0 m(k) \leq m(k-1) \leq \beta_1 m(k)$$

holds with $1 < \beta_0 \leq \beta_1$.

Typically β_1 is chosen so that the total work required for a multigrid cycle is no greater than the work required for application of the stiffness matrix on the finest level. This condition is satisfied, if for instance, $m(k) = 2^{J-k}$.

The following theorem is the main result of this paper.

THEOREM 4.1. *Assume that (A.1) and (A.2) hold. There exists an α and $M > 0$ independent of J such that*

$$\eta^{-1} \tilde{A}(u, u) \leq \tilde{A}(B_J A_J u, u) \leq \eta \tilde{A}(u, u) \quad \text{for all } u \in M_J$$

with $\eta = \frac{M+m(J)^{\beta/2}}{m(J)^{\beta/2}}$.

The theorem shows that B_J is a uniform preconditioner for the linear system arising from mortar finite element discretization using M_J even if $m(J) = 1$. Increasing $m(J)$ gives a somewhat better rate of convergence but increases the cost of applying B_J . It suffices to choose α above so that $\alpha < 1/C_1$ where C_1 is as in Lemma 4.5.

We use the following lemmas to prove Theorem 4.1. Their proofs will be given in section 5. First, we state a lemma that is a consequence of regularity which will be used in the proof of Condition (C.2).

LEMMA 4.2. *If (A.1) holds, then*

$$\|(I - I_k P_{k-1})u\| \leq Ch_k^\beta \|A_k u\|_{0,\Omega}^\beta \tilde{A}(u, u)^{(1-\beta)/2}$$

holds for all u in M_k .

The next three lemmas are useful in analyzing the smoothing operator. We begin with a lemma from the theory of additive preconditioners.

LEMMA 4.3. *Let the space V be a sum of subspaces $\sum_{i=1}^l V_i$. For $i = 1, 2, \dots, l$, let B_i be a symmetric positive definite operator on V_i and Q_i be the L^2 projection onto V_i . Then for $B = \sum_{i=1}^l B_i Q_i$,*

$$(B^{-1}u, u) = \inf_{\substack{u_i \in V_i \\ u = \sum_{i=1}^l u_i}} \left(\sum_{i=1}^l (B_i^{-1}u_i, u_i) \right)$$

holds for all u in V .

Lemma 4.3 may be found in [9] or stated in a different form in [17, Chapter 4]. We do not prove it here. The following two lemmas are used in the proof of Condition (C.1).

LEMMA 4.4. *For R_k defined by (4.1), there exists a constant $C_R = C_R(\alpha)$ independent of k such that (4.2) holds for all u in M_k .*

LEMMA 4.5. *For all u in M_k , there is a number C_1 not depending on J such that*

$$(A_k u, u) \leq C_1 \sum_{i=1}^{N_k} c_i^2 \tilde{A}(\phi_k^i, \phi_k^i),$$

where $u = \sum_{i=1}^{N_k} c_i \phi_k^i$ is the nodal basis decomposition.

We now prove the theorem.

Proof of Theorem 4.1. We apply the theorem for variable V-cycle in [6, Theorem 4.6]. This requires verification of Conditions (C.1) and (C.2).

Because of Lemma 4.4, (C.1) follows if we show that $I - R_k A_k$ is nonnegative, i.e., for all $u \in M_k$,

$$(A_k R_k A_k u, u) \leq (A_k u, u).$$

This is equivalent to showing that for all $u \in M_k$,

$$(A_k u, u) \leq (R_k^{-1} u, u).$$

Fix $u \in M_k$ and let $u = \sum_{i=1}^{N_k} c_i \phi_k^i$ be its nodal basis decomposition. Applying Lemma 4.3 gives

$$(R_k^{-1} u, u) = \frac{1}{\alpha} \sum_{i=1}^{N_k} (A_{k,i} c_i \phi_k^i, c_i \phi_k^i) = \frac{1}{\alpha} \sum_{i=1}^{N_k} c_i^2 \tilde{A}(\phi_k^i, \phi_k^i).$$

The nonnegativity of $I - R_k A_k$ follows provided that α is taken to be less than or equal to $1/C_1$ where C_1 is as in Lemma 4.5.

Condition (C.2) is immediately seen to hold from Lemma 4.2. Indeed,

$$\begin{aligned} \tilde{A}((I - I_k P_{k-1})u, u) &\leq C \|(I - I_k P_{k-1})u\| \|u\| \\ &\leq C \left(\frac{\|A_k u\|_{0,\Omega}^2}{\lambda_k} \right)^{\beta/2} \tilde{A}(u, u)^{1-\beta/2}. \end{aligned}$$

Here we have used the fact that $\lambda_k \leq C h_k^{-2}$. This proves (C.2) and thus completes the proof of the theorem. \square

5. Proof of the lemmas. As a first step in proving Lemma 4.2, we prove that the operators $\{I_k\}$ approximate the identity operator I . After proving Lemma 4.2, we state and prove two lemmas used in the proof of Lemmas 4.4 and 4.5.

LEMMA 5.1. For any $u \in M_{k-1}$,

$$(5.1) \quad \|(I - I_k)u\| \leq C \inf_{v \in \mathcal{V}} \|u - v\| \quad \text{and}$$

$$(5.2) \quad \|(I - I_k)u\|_{0,\Omega} \leq Ch_k \|u\|.$$

Proof. Fix $u \in M_{k-1}$. By definition, $(I_k - I)u = \sum_{\gamma \in Z} \mathcal{E}_{k,\gamma}u$. Since $\mathcal{E}_{k,\gamma}u$ is zero on every interior vertex of the mesh in $\Omega_{NM}(\gamma)$,

$$(5.3) \quad \|\mathcal{E}_{k,\gamma}u\|_{1,\Omega_{NM}(\gamma)}^2 \approx \sum_{y_k^i} (\mathcal{E}_{k,\gamma}u)(y_k^i)^2 \approx h_k^{-1} \|\mathcal{E}_{k,\gamma}u\|_{0,\gamma}^2.$$

The above sum is taken over the vertices y_k^i of the $\Omega_{NM}(\gamma)$ mesh that lie on γ . Here and elsewhere \approx denotes equivalence with constants independent of h_k and $\|\mathcal{E}_{k,\gamma}u\|_{0,\gamma}$ denotes the $L^2(\gamma)$ norm of the nonmortar trace of $\mathcal{E}_{k,\gamma}u$. By the L^2 stability of $\Pi_{k,\gamma}$,

$$(5.4) \quad \begin{aligned} \|\mathcal{E}_{k,\gamma}u\|_{0,\gamma}^2 &= \|\Pi_{k,\gamma}(u_\gamma^M - u_\gamma^{NM})\|_{0,\gamma}^2 \\ &\leq C \|u_\gamma^M - u_\gamma^{NM}\|_{0,\gamma}^2. \end{aligned}$$

Since u is in M_{k-1} , denoting $u_\gamma^M - u_\gamma^{NM}$ by e , we have that for any $\nu \in S_{k-1}(\gamma)$,

$$(e, e)_\gamma = (e, e - \nu)_\gamma \leq \|e\|_{0,\gamma} \|e - \nu\|_{0,\gamma},$$

where $(\cdot, \cdot)_\gamma$ denotes the $L^2(\gamma)$ inner-product. It then follows that

$$\|\mathcal{E}_{k,\gamma}u\|_{0,\gamma} \leq \inf_{\nu \in S_{k-1}(\gamma)} \|e - \nu\|_{0,\gamma}.$$

If q denotes the L^2 orthogonal projection into $S_{k-1}(\gamma)$, then by the approximation properties of $S_{k-1}(\gamma)$, we have $\|e - q(e)\|_{0,\gamma} \leq Ch_k \|e\|_{1,\gamma}$. Trivially, we also have that $\|e - q(e)\|_{0,\gamma} \leq \|e\|_{0,\gamma}$. Thus the theory of interpolation of operators gives that

$$(5.5) \quad \|\mathcal{E}_{k,\gamma}u\|_{0,\gamma} \leq Ch_k^{1/2} \|u_\gamma^M - u_\gamma^{NM}\|_{1/2,\gamma}.$$

Since $u_\gamma^M - u_\gamma^{NM}$ is equal to $(u - v)_\gamma^M - (u - v)_\gamma^{NM}$ whenever v is in \mathcal{V} , the first inequality of the lemma readily follows when (5.5) is combined with (5.3), and a trace theorem is applied.

To prove the last inequality of the lemma, we start with

$$\|(I - I_k)u\|_{0,\Omega}^2 \leq Ch_k \sum_{\gamma \in Z} \|\mathcal{E}_{k,\gamma}u\|_{0,\gamma}^2.$$

Applying (5.5), we get

$$\|(I - I_k)u\|_{0,\Omega}^2 \leq Ch_k^2 \sum_{\gamma \in Z} \|u_\gamma^M - u_\gamma^{NM}\|_{1/2,\gamma}^2,$$

from which (5.2) follows after using the triangle inequality and a trace theorem. □

Proof of Lemma 4.2. The proof is broken into two parts. First, we prove that

$$(5.6) \quad \|(I - I_k P_{k-1})u\| \leq C(h_k^\beta \|A_k u\|_{-1+\beta, \Omega} + h_k \|A_k u\|_{0, \Omega})$$

holds for all u in M_k . Next, we show that

$$(5.7) \quad \|A_k u\|_{-1+\beta, \Omega} \leq C\tilde{A}(u, u)^{(1-\beta)/2} \|A_k u\|_{0, \Omega}^\beta$$

holds for all u in M_k . Clearly the lemma follows using (5.7) to bound the first term on the right-hand side of (5.6) and the fact that $\lambda_k \leq Ch_k^{-2}$.

Fix u in M_k and set $g = A_k u$. Then u solves

$$\tilde{A}(u, \phi) = (g, \phi) \quad \text{for all } \phi \in M_k.$$

Let $w \in \mathcal{V}$ be the solution of

$$(5.8) \quad A(w, \phi) = (g, \phi) \quad \text{for all } \phi \in \mathcal{V}.$$

Now u is the mortar finite element approximation to w from M_k and hence by (3.4),

$$(5.9) \quad \|u - w\| \leq Ch_k^\beta \|w\|_{1+\beta, \Omega}.$$

By the triangle inequality,

$$(5.10) \quad \|u - I_k P_{k-1} u\| \leq Ch_k^\beta \|w\|_{1+\beta, \Omega} + \|w - I_k P_{k-1} u\|.$$

To estimate the second term of (5.10), we start by writing $P_{k-1} u = v_1 + v_2$ where $v_1 \in M_{k-1}$ solves

$$\tilde{A}(v_1, \phi) = (g, \phi) \quad \text{for all } \phi \in M_{k-1}.$$

The remainder v_2 satisfies

$$(5.11) \quad \tilde{A}(v_2, \phi) = (g, (I_k - I)\phi) \quad \text{for all } \phi \in M_{k-1}.$$

Then

$$\begin{aligned} \|w - I_k P_{k-1} u\| &\leq \|w - v_1\| + \|I_k v_2\| + \|(I - I_k)v_1\| \\ &\leq C \|w - v_1\| + C \|v_2\| \\ &\leq Ch_k^\beta \|w\|_{1+\beta, \Omega} + C \|v_2\|, \end{aligned}$$

where we have used Lemma 5.1 and (3.4). Using the coercivity of the form $\tilde{A}(\cdot, \cdot)$, we find from (5.11) that

$$\begin{aligned} \|v_2\|^2 &\leq C\tilde{A}(v_2, v_2) = C(A_k u, (I_k - I)v_2) \\ &\leq \|A_k u\|_{0, \Omega} \|(I - I_k)v_2\|_{0, \Omega}. \end{aligned}$$

The last inequality of Lemma 5.1 now yields that $\|v_2\| \leq Ch_k \|A_k u\|_{0, \Omega}$. Combining these estimates we get that

$$\|w - I_k P_{k-1} u\| \leq Ch_k^\beta \|w\|_{1+\beta, \Omega} + Ch_k \|A_k u\|_{0, \Omega}.$$

Using this in (5.10) and applying Assumption (A.1) proves (5.6).

We next prove (5.7). Fix u in M_k . Since $\|\cdot\|_{-1+\beta,\Omega}$ is the norm on the space in the interpolation scale between \mathcal{V}' and $L^2(\Omega)$,

$$\|A_k u\|_{-1+\beta,\Omega} \leq \|A_k u\|_{-1,\Omega}^{1-\beta} \|A_k u\|_{0,\Omega}^\beta.$$

Thus, it suffices to prove that

$$(5.12) \quad \|A_k u\|_{-1,\Omega} \leq C \tilde{A}(u, u)^{1/2}.$$

Given ψ in \mathcal{V} , we will construct $\psi_k = \psi_k(\psi) \in M_k$ satisfying

$$(5.13) \quad \|\psi_k\| \leq C \|\psi\|_{1,\Omega} \quad \text{and}$$

$$(5.14) \quad \|\psi - \psi_k\|_{0,\Omega} \leq C h_k \|\psi\|_{1,\Omega}.$$

Assuming such a ψ_k exists, we have

$$\begin{aligned} \|A_k u\|_{-1,\Omega} &= \sup_{\psi \in \mathcal{V}} \frac{(A_k u, \psi)}{\|\psi\|_{1,\Omega}} \\ &\leq \sup_{\psi \in \mathcal{V}} \frac{(A_k u, \psi - \psi_k)}{\|\psi\|_{1,\Omega}} + \sup_{\psi \in \mathcal{V}} \frac{(A_k u, \psi_k)}{\|\psi\|_{1,\Omega}}. \end{aligned}$$

Inequality (5.12) then follows from

$$\begin{aligned} \|A_k u\|_{-1,\Omega} &\leq \sup_{\psi \in \mathcal{V}} \frac{\|A_k u\|_{0,\Omega} \|\psi - \psi_k\|_{0,\Omega}}{\|\psi\|_{1,\Omega}} + \sup_{\psi \in \mathcal{V}} \frac{\tilde{A}(u, \psi_k)}{\|\psi\|_{1,\Omega}} \\ &\leq C \left(h_k \|A_k u\|_{0,\Omega} + \tilde{A}(u, u)^{1/2} \sup_{\psi \in \mathcal{V}} \frac{\|\psi_k\|_{1,\Omega}}{\|\psi\|_{1,\Omega}} \right) \\ &\leq C h_k \lambda_k^{1/2} \tilde{A}(u, u)^{1/2} + C \tilde{A}(u, u)^{1/2} \\ &\leq C \tilde{A}(u, u)^{1/2}. \end{aligned}$$

To complete the proof, we need only construct ψ_k satisfying (5.13) and (5.14). For $\psi \in \mathcal{V}$, let $\tilde{\psi}_k \in \tilde{M}_k$ be the L^2 projection of ψ into \tilde{M}_k . This projection is local on Ω_i and satisfies (see [8])

$$(5.15) \quad \|\tilde{\psi}_k\| \leq C \|\psi\|_{1,\Omega} \quad \text{and}$$

$$(5.16) \quad \|\psi - \tilde{\psi}_k\|_{0,\Omega} \leq C h_k \|\psi\|_{1,\Omega}.$$

To construct ψ_k , we modify $\tilde{\psi}_k$ so that the result is in M_k , i.e.,

$$\psi_k = \tilde{\psi}_k + \sum_{\gamma \in Z} \mathcal{E}_{k,\gamma}(\tilde{\psi}_k).$$

We will now show that ψ_k defined above satisfies (5.13). We start with

$$\|\psi_k\| \leq \|\tilde{\psi}_k\| + \sum_{\gamma \in Z} \|\mathcal{E}_{k,\gamma} \tilde{\psi}_k\|.$$

Using (5.15) on the first term on right-hand side and using (5.3) on the remaining, we get

$$(5.17) \quad \|\psi_k\|^2 \leq C \left(\|\psi\|_{1,\Omega}^2 + h_k^{-1} \sum_{\gamma \in Z} \|\mathcal{E}_{k,\gamma} \tilde{\psi}_k\|_{0,\gamma}^2 \right).$$

Note that $\|\mathcal{E}_{k,\gamma} \tilde{\psi}_k\|_{0,\gamma} \leq C \|(\tilde{\psi}_k)_\gamma^M - (\tilde{\psi}_k)_\gamma^{NM}\|_{0,\gamma}$ by (5.4). Since ψ is in $H^1(\Omega)$, its trace on γ is in $L^2(\gamma)$. Moreover, ψ_γ^M and ψ_γ^{NM} are equal. Hence,

$$\begin{aligned} \|\mathcal{E}_{k,\gamma} \tilde{\psi}_k\|_{0,\gamma} &\leq C \left\| (\tilde{\psi}_k - \psi)_\gamma^M - (\tilde{\psi}_k - \psi)_\gamma^{NM} \right\|_{0,\gamma} \\ &\leq C \left\| \tilde{\psi}_k - \psi \right\|_{0,\Omega_M(\gamma)}^{1/2} \left\| \tilde{\psi}_k - \psi \right\|_{1,\Omega_M(\gamma)}^{1/2} \\ &\quad + C \left\| \tilde{\psi}_k - \psi \right\|_{0,\Omega_{NM}(\gamma)}^{1/2} \left\| \tilde{\psi}_k - \psi \right\|_{1,\Omega_{NM}(\gamma)}^{1/2}, \end{aligned}$$

where in the last step we have used a trace inequality. Using (5.15) and (5.16), we then have

$$(5.18) \quad \|\mathcal{E}_{k,\gamma} \tilde{\psi}_k\|_{0,\gamma} \leq Ch_k^{1/2} \left(\|\psi\|_{1,\Omega_M(\gamma)} + \|\psi\|_{1,\Omega_{NM}(\gamma)} \right).$$

Combining (5.18) and (5.17) gives (5.13).

It now remains only to prove (5.14). By the triangle inequality,

$$\|\psi - \psi_k\|_{0,\Omega} \leq \left\| \psi - \tilde{\psi}_k \right\|_{0,\Omega} + \left\| \tilde{\psi}_k - \psi_k \right\|_{0,\Omega}.$$

The first term on the right-hand side is readily bounded as required by (5.16). For the second term, as in (5.3),

$$\left\| \tilde{\psi}_k - \psi_k \right\|_{0,\Omega} \leq Ch_k^{1/2} \sum_{\gamma \in Z} \|\mathcal{E}_{k,\gamma} \tilde{\psi}_k\|_{0,\gamma}.$$

Inequality (5.14) now follows immediately from (5.18). This completes the proof of Lemma 4.2. \square

We are left to prove the lemmas involving the smoother R_k . A critical ingredient in this analysis involves the decay properties of the projector $\Pi_{k,\gamma}$ away from the support of the data. Specifically, we use the following lemma.

LEMMA 5.2. *Let the support of $v \in L^2(\gamma)$ be contained in a connected subset $\sigma \subseteq \gamma$. Then there is a constant c such that for any set $\kappa \subseteq \gamma$ disjoint from σ ,*

$$\|\Pi_{k,\gamma} v\|_{0,\kappa} \leq C \exp\left(-c \frac{\text{dist}(\kappa, \sigma)}{h_k}\right) \|v\|_{0,\gamma},$$

where $\text{dist}(\kappa, \sigma)$ is the distance between the sets κ and σ .

Remark 5.3. Estimates similar to those in the above lemma for the L^2 -orthogonal projection were given by Descloux [12]. Note that $\Pi_{k,\gamma}$ is not an L^2 -orthogonal projection. For completeness, we include a proof for our case which is a modification of one given in [21, Chapter 5].

Proof. Recall that a $\gamma \in Z$ is partitioned into subintervals $\omega_{k,i}$ by the vertices $x_{k,\gamma}^i$, $i = 0, \dots, N$ of the mesh on $\Omega_{NM}(\gamma)$. Define the set r_0 as the union of those

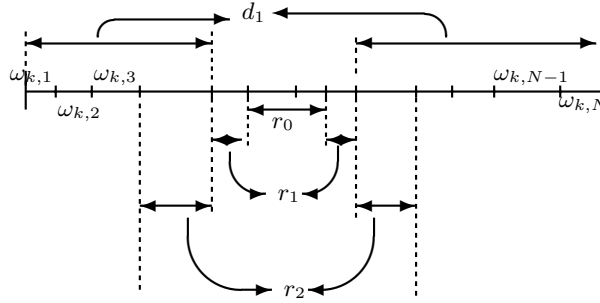


FIG. 5.1. An interface segment.

subintervals which intersect the support of v . Following the presentation in [21], define r_j , $j = 1, 2, \dots$, recursively, by letting r_m be the union of those subintervals of γ that are not in $\cup_{l < m} r_l$ and which are neighbors of the subintervals of this set (see Figure 5.1). Further, let $d_m = \cup_{l > m} r_l$.

We will now show that the L^2 norm of $\Pi_{k,\gamma}v$ on d_m can be bounded by a constant times its L^2 norm on r_m . For all $\chi \in S_k(\gamma)$ with support of χ disjoint from r_0 , we have

$$(5.19) \quad (\Pi_{k,\gamma}v, \chi) = (v, \chi) = 0.$$

Let $\chi_m \in S_k(\gamma)$, for $m \geq 1$, be defined by

$$\chi_m(x_{k,\gamma}^j) = \begin{cases} \Pi_{k,\gamma}v(x_{k,\gamma}^j) & \text{for } x_{k,\gamma}^j \in \bar{d}_m, \\ 0 & \text{otherwise,} \end{cases}$$

for $j = 1, \dots, N - 1$. Let $\varepsilon = \omega_{k,1} \cup \omega_{k,N}$. Clearly, (5.19) holds with χ_m in place of χ . Moreover, $\chi_m(x) = \Pi_{k,\gamma}v(x)$ for $x \in d_m \setminus \varepsilon$, and it vanishes on $\gamma \setminus d_{m-1}$. Then

$$0 = (\chi_m, \Pi_{k,\gamma}v) = \int_{d_m \setminus \varepsilon} |\Pi_{k,\gamma}v|^2 ds + \int_{d_m \cap \varepsilon} \chi_m \Pi_{k,\gamma}v ds + \int_{r_m} \chi_m \Pi_{k,\gamma}v ds.$$

Note that on each subinterval of $d_m \cap \varepsilon$, χ_m is constant, and it takes the value of $\Pi_{k,\gamma}v$ at the interior endpoint. Also, on the subintervals of r_m , χ_m is either identically zero (if that subinterval is part of $r_m \cap \varepsilon$) or takes the value of $\Pi_{k,\gamma}v$ on one endpoint and zero on the other endpoint. From these observations, it is easy to conclude that

$$\int_{d_m \cap \varepsilon} \chi_m \Pi_{k,\gamma}v ds \geq C \|\Pi_{k,\gamma}v\|_{0,\varepsilon \cap d_m}^2$$

and

$$\int_{r_m} |\Pi_{k,\gamma}v| |\chi_m| ds \leq C \|\Pi_{k,\gamma}v\|_{0,r_m}^2.$$

Thus,

$$\begin{aligned} C \|\Pi_{k,\gamma}v\|_{0,d_m}^2 &\leq \int_{d_m \setminus \varepsilon} |\Pi_{k,\gamma}v|^2 ds + \int_{d_m \cap \varepsilon} \chi_m \Pi_{k,\gamma}v ds \\ &= - \int_{r_m} \chi_m \Pi_{k,\gamma}v ds \leq C \|\Pi_{k,\gamma}v\|_{0,r_m}^2. \end{aligned}$$

Letting $q_m = \|\Pi_{k,\gamma} v\|_{0,d_m}^2$, the above inequality can be rewritten as $q_m \leq C(q_{m-1} - q_m)$. It immediately follows that

$$q_m \leq \frac{C}{1+C} q_{m-1} \leq \dots \leq \left(\frac{C}{1+C}\right)^m \|\Pi_{k,\gamma} v\|_{0,\gamma}^2.$$

The lemma easily follows from (3.6) and the observation that the distance between κ and σ is $O(mh)$. \square

Proof of Lemma 4.4. Fix $u \in M_k$ and let $u = \sum_{i=1}^{N_k} c_i \phi_k^i$ be the nodal basis decomposition. By Lemma 4.3,

$$\begin{aligned} (R_k^{-1}u, u) &= \frac{1}{\alpha} \sum_{i=1}^l (A_{k,i}(c_i \phi_k^i), c_i \phi_k^i) \\ &\leq \frac{\lambda_k}{\alpha} \sum_{i=1}^l c_i^2 (\phi_k^i, \phi_k^i). \end{aligned}$$

Note that the L^2 norm of every basis function ϕ_k^i is $O(h_k^2)$. Indeed, this is a standard estimate for those basis functions that coincide with a usual finite element nodal basis function on a subdomain. For the remaining basis functions, this follows from the exponential decay given by Lemma 5.2. Thus,

$$(5.20) \quad (R_k^{-1}u, u) \leq \frac{C\lambda_k h_k^2}{\alpha} \sum_{i=1}^{N_k} c_i^2.$$

On each subdomain Ω_j we have that

$$\|u\|_{0,\Omega_j}^2 \approx h_k^2 \left(\sum_{i=1}^{\tilde{N}_k} u(y_k^i)^2 \right).$$

Combining the above inequalities gives

$$(R_k^{-1}u, u) \leq \frac{C\lambda_k}{\alpha} \|u\|_{0,\Omega}^2.$$

The above inequality is equivalent to (4.2) and thus completes the proof of the lemma. \square

The proof of Lemma 4.5 requires a strengthened Cauchy–Schwarz inequality which we provide in the next lemma. First, we introduce some notation. Define the index sets $\tilde{\mathcal{N}}_k^\gamma$ and \mathcal{N}_k^γ by

$$\begin{aligned} \tilde{\mathcal{N}}_k^\gamma &= \{i : y_k^i \in \gamma \cap \bar{\Omega}_{NM(\gamma)}\}, \\ \mathcal{N}_k^\gamma &= \{i : y_k^i \in \bar{\gamma} \text{ and } i \notin \tilde{\mathcal{N}}_k^\gamma\}. \end{aligned}$$

Also denote the set $\cup\{\mathcal{N}_k^\gamma : \gamma \in Z\}$ by \mathcal{N}_k^Γ .

LEMMA 5.4. *Let ϕ_k^i and ϕ_k^j be two basis functions of M_k with $i, j \in \mathcal{N}_k^\Gamma$. Let y_k^i and y_k^j be the corresponding vertices. Then $\tilde{A}(\phi_k^i, \phi_k^j)$ satisfies*

$$\tilde{A}(\phi_k^i, \phi_k^j) \leq C \exp\left(-c \frac{|y_k^i - y_k^j|}{h_k}\right) \tilde{A}(\phi_k^i, \phi_k^i)^{1/2} \tilde{A}(\phi_k^j, \phi_k^j)^{1/2},$$

where C and c are constants independent of k .

Proof. First, consider the case when y_k^i and y_k^j are on a same open interface segment $\gamma \in Z$. Let Δ_M denote the set of triangles that have at least one vertex on $\bar{\gamma}$ and are contained in $\Omega_{M(\gamma)}$. Similarly let Δ_{NM} denote the set of triangles that have at least one vertex on $\bar{\gamma}$ and are contained in $\Omega_{NM(\gamma)}$. Since ϕ_k^i and ϕ_k^j are simultaneously nonzero only on triangles in Δ_M or Δ_{NM} , we have

$$(5.21) \quad \tilde{A}(\phi_k^i, \phi_k^j) = \sum_{\tau \in \Delta_M} A_\tau(\phi_k^i, \phi_k^j) + \sum_{\tau \in \Delta_{NM}} A_\tau(\phi_k^i, \phi_k^j),$$

where $A_\tau(u, v) = \int_\tau \nabla u \cdot \nabla v \, dx$. The first sum obviously satisfies the required inequality, because this sum is zero whenever y_k^i and y_k^j are not vertices of the same triangle in Δ_M .

Now consider a triangle $\tau \in \Delta_{NM}$. Recall that γ was subdivided by the nonmortar mesh into subintervals $\omega_{k,i}$, $i = 1, \dots, N$. Let ω_τ denote the union of two or more of these subintervals which have the vertices of τ as an endpoint (see Figure 5.2). Then because ϕ_k^i and ϕ_k^j are zero at the vertices of τ that are not in ω_τ ,

$$(5.22) \quad A_\tau(\phi_k^i, \phi_k^j) \leq |\phi_k^i|_{1,\tau} |\phi_k^j|_{1,\tau} \leq h_k^{-1} \|(\phi_k^i)_\gamma^{NM}\|_{0,\omega_\tau} \|(\phi_k^j)_\gamma^{NM}\|_{0,\omega_\tau}.$$

Now, recall that ϕ_k^i and ϕ_k^j are obtained from $\tilde{\phi}_k^i$ and $\tilde{\phi}_k^j$, respectively, as described by (3.9). Denote by s_i and s_j the supports of the nonzero traces $\tilde{\phi}_k^i|_\gamma$ and $\tilde{\phi}_k^j|_\gamma$, respectively. Then by Lemma 5.2,

$$\begin{aligned} A_\tau(\phi_k^i, \phi_k^j) &\leq C h_k^{-1} \exp\left(-\frac{c}{h_k} [\text{dist}(s_i, \omega_\tau) + \text{dist}(s_j, \omega_\tau)]\right) \|\tilde{\phi}_k^i\|_{0,\gamma} \|\tilde{\phi}_k^j\|_{0,\gamma} \\ &\leq C \exp\left(-\frac{c}{h_k} [\text{dist}(s_i, \omega_\tau) + \text{dist}(s_j, \omega_\tau)]\right) \|\tilde{\phi}_k^i\|_{1,\Omega_{M(\gamma)}} \|\tilde{\phi}_k^j\|_{1,\Omega_{M(\gamma)}}. \end{aligned}$$

Now, if $|\omega_\tau|$ denotes the length of ω_τ , it may easily be seen that

$$\text{dist}(s_i, \omega_\tau) + \text{dist}(s_j, \omega_\tau) + |\omega_\tau| \geq \text{dist}(s_i, s_j).$$

Further, by quasi uniformity,

$$\text{dist}(s_i, s_j) \geq |y_k^i - y_k^j| - Ch_k.$$

Split the sum over $\tau \in \Delta_{NM}$ in (5.21) into a sum over triangles which have a vertex lying in between y_k^i and y_k^j on γ , and a sum over the remaining triangles in Δ_{NM} . We denote the former set of triangles as Δ_{NM}^{in} and the latter as Δ_{NM}^{out} . Note that the number of triangles in Δ_{NM}^{in} is bounded by $C|y_k^i - y_k^j|/h_k$.

We first consider triangles in Δ_{NM}^{in} . The observations of the previous paragraph yield

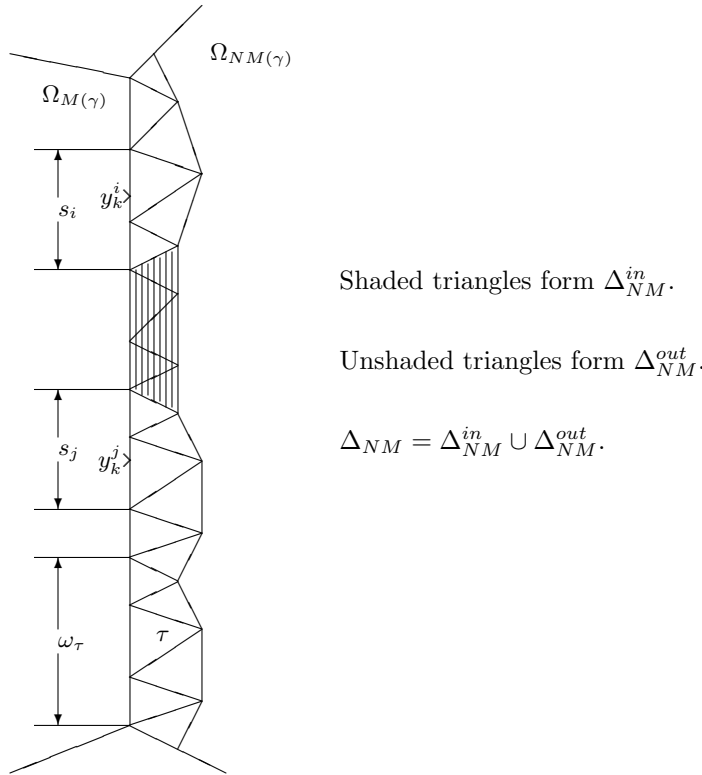


FIG. 5.2. Illustrating the notations in the proof of Lemma 5.4.

$$\begin{aligned}
 & \sum_{\tau \in \Delta_{NM}^{in}} A_{\tau}(\phi_k^i, \phi_k^j) \\
 & \leq C \exp\left(-\frac{c}{h_k} \text{dist}(s_i, s_j)\right) \|\tilde{\phi}_k^i\|_{1, \Omega_M(\gamma)} \|\tilde{\phi}_k^j\|_{1, \Omega_M(\gamma)} \left(\sum_{\tau \in \Delta_{NM}^{in}} 1\right) \\
 & \leq C \frac{|y_k^i - y_k^j|}{h_k} \exp\left(-c \frac{|y_k^i - y_k^j|}{h_k}\right) \|\tilde{\phi}_k^i\|_{1, \Omega_M(\gamma)} \|\tilde{\phi}_k^j\|_{1, \Omega_M(\gamma)} \\
 (5.23) \quad & \leq C \exp\left(-c \frac{|y_k^i - y_k^j|}{2h_k}\right) \|\tilde{\phi}_k^i\|_{1, \Omega_M(\gamma)} \|\tilde{\phi}_k^j\|_{1, \Omega_M(\gamma)}.
 \end{aligned}$$

Now, for the sum over triangles in Δ_{NM}^{out} , observe that one of the distances, $\text{dist}(\omega_{\tau}, s_i)$ or $\text{dist}(\omega_{\tau}, s_j)$, is greater than $\text{dist}(s_i, s_j)$. Hence

$$\begin{aligned}
 \sum_{\tau \in \Delta_{NM}^{out}} A_{\tau}(\phi_k^i, \phi_k^j) & \leq C \exp\left(-c \frac{|y_k^i - y_k^j|}{h_k}\right) \|\tilde{\phi}_k^i\|_{1, \Omega_M(\gamma)} \|\tilde{\phi}_k^j\|_{1, \Omega_M(\gamma)} \\
 & \quad \times \sum_{\tau \in \Delta_{NM}^{out}} \exp\left(-c \frac{\text{dist}(\omega_{\tau}, s_i \cup s_j)}{h_k}\right).
 \end{aligned}$$

Since the sum on the right-hand side can be bounded by a summable geometric series,

we have

$$(5.24) \quad \sum_{\tau \in \Delta_{NM}^{out}} A_\tau(\phi_k^i, \phi_k^j) \leq C \exp\left(-c \frac{|y_k^i - y_k^j|}{h_k}\right) \|\tilde{\phi}_k^i\|_{1, \Omega_M(\gamma)} \|\tilde{\phi}_k^j\|_{1, \Omega_M(\gamma)}.$$

Thus, (5.23), (5.24), and (5.21) give

$$\tilde{A}(\phi_k^i, \phi_k^j) \leq C \exp\left(-c \frac{|y_k^i - y_k^j|}{2h_k}\right) \|\phi_k^i\| \|\phi_k^j\|.$$

This, along with the coercivity of $\tilde{A}(\cdot, \cdot)$ on $M_k \times M_k$, proves the lemma when y_k^i and y_k^j lie on the same γ .

Now suppose one of y_k^i and y_k^j , say y_k^i , is an endpoint of $\bar{\gamma}$. Then (5.22) holds for all $\tau \in \Delta_{NM}$, except possibly a triangle at all of whose vertices ϕ_k^i is nonzero. This triangle must have y_k^i as one of its vertices. If τ is this triangle, then clearly

$$A_\tau(\phi_k^i, \phi_k^j) \leq |\phi_k^i|_{1, \tau} h_k^{-1/2} |\phi_k^j|_{0, \omega_\tau} \leq C \|\phi_k^i\|_{1, \tau} \exp\left(-\frac{c}{h_k} \text{dist}(s_i, s_j)\right) \|\tilde{\phi}_k^j\|_{1, \Omega_M(\gamma)}.$$

Hence, all the arguments for the previous case hold for this case too. The case when y_k^i is one endpoint of γ and y_k^j is the other is handled in a similar fashion by considering separately the two possible triangles where (5.22) may not hold.

To conclude the proof, it now suffices to consider the case when $y_k^i \in \gamma_1$ and $y_k^j \in \gamma_2$ with $\gamma_1 \neq \gamma_2$, and $\gamma_1, \gamma_2 \in Z$. Then $\tilde{A}(\phi_k^i, \phi_k^j)$ is zero unless there is a triangle τ in \mathcal{T}_k which has one of its vertices contained in γ_1 and another contained in γ_2 . It is easy to see that the required estimate follows in the latter case too. \square

Proof of Lemma 4.5. Split u into a function u_0 that vanishes on the interface Γ and a function u_Γ that is a linear combination of ϕ_k^i , with $i \in \mathcal{N}_k^\Gamma$. By the triangle inequality,

$$\tilde{A}(u, u) \leq 2[\tilde{A}(u_0, u_0) + \tilde{A}(u_\Gamma, u_\Gamma)].$$

On each triangle τ in \mathcal{T}_k ,

$$u_0 = \sum_{j=1}^3 c_{i(\tau;j)} \phi_k^{i(\tau;j)} \text{ on } \tau,$$

where $i(\tau; j)$, $j = 1, 2, 3$, are the vertices of τ . Applying the arithmetic-geometric mean inequality gives

$$\begin{aligned} \tilde{A}(u_0, u_0) &= \sum_{\tau \in \mathcal{T}_k} A_\tau(u_0, u_0) \\ &\leq \sum_{\tau \in \mathcal{T}_k} 3 \sum_{j=1}^3 c_{i(\tau;j)}^2 A_\tau(\phi_k^{i(\tau;j)}, \phi_k^{i(\tau;j)}) \\ &= 3 \sum_{i \notin \mathcal{N}_k^\Gamma} c_i^2 \sum_{\tau \in \mathcal{T}_k} A_\tau(\phi_k^i, \phi_k^i) \\ &= 3 \sum_{i \notin \mathcal{N}_k^\Gamma} c_i^2 \tilde{A}(\phi_k^i, \phi_k^i). \end{aligned}$$

All that remains is to estimate $\tilde{A}(u_\Gamma, u_\Gamma)$. We clearly have

$$\tilde{A}(u_\Gamma, u_\Gamma) = \sum_{i,j \in \mathcal{N}_\Gamma} c_i c_j \tilde{A}(\phi_k^i, \phi_k^j).$$

Applying Lemma 5.4 gives

$$\begin{aligned} \tilde{A}(u_\Gamma, u_\Gamma) &\leq C \sum_{i,j \in \mathcal{N}_k^\Gamma} c_i c_j \exp\left(-c \frac{|y_k^i - y_k^j|}{h_k}\right) \tilde{A}(\phi_k^i, \phi_k^i)^{1/2} \tilde{A}(\phi_k^j, \phi_k^j)^{1/2} \\ &\leq C \|\mathbf{M}\|_{\ell^2} \sum_{i \in \mathcal{N}_k^\Gamma} c_i^2 \tilde{A}(\phi_k^i, \phi_k^i). \end{aligned}$$

Here \mathbf{M} is the matrix with entries

$$\mathbf{M}_{ij} = \exp\left(-c \frac{|y_k^i - y_k^j|}{h_k}\right)$$

and

$$\|\mathbf{M}\|_{\ell^2} = \sup_{\zeta \in \mathbb{R}^{|\mathcal{N}_k^\Gamma|}} \frac{(\mathbf{M}\zeta) \cdot \zeta}{\zeta \cdot \zeta},$$

where $|\mathcal{N}_k^\Gamma|$ denotes the cardinality of \mathcal{N}_k^Γ and \cdot indicates the standard dot product in $\mathbb{R}^{|\mathcal{N}_k^\Gamma|}$.

To conclude the proof, it suffices to show that $\|\mathbf{M}\|_{\ell^2}$ is bounded by a constant independent of h_k . Note that $\|\mathbf{M}\|_{\ell^2}$ is equal to the spectral radius of \mathbf{M} and consequently, can be bounded by any induced norm. So,

$$\|\mathbf{M}\|_{\ell^2} \leq \max_{i \in \mathcal{N}_k^\Gamma} \sum_{j \in \mathcal{N}_k^\Gamma} \mathbf{M}_{ij}.$$

For every fixed i , the sum on the right-hand side can be enlarged to run over all vertices of the mesh \mathcal{T}_k , and then one obtains

$$\sum_{j \in \mathcal{N}_k^\Gamma} \mathbf{M}_{ij} \leq \sum_{y_k^j \in \mathcal{T}_k} \exp\left(-c \frac{|y_k^j - y_k^i|}{h_k}\right) \leq C \iint_{\mathbb{R}^2} \exp(-c|y|) dy.$$

Thus, $\|\mathbf{M}\|_{\ell^2} \leq C$. \square

6. Implementation. This section will describe some details of implementing the mortar method and the preconditioner B_J . Since we shall be using a preconditioned iteration, all that is necessary is the implementation of the action of the stiffness matrix and that of the preconditioner.

Let \mathbf{A}_k denote the stiffness matrix for the mortar finite element method, i.e., $[\mathbf{A}_k]_{ij} = \tilde{A}(\phi_k^j, \phi_k^i)$. Let

$$(6.1) \quad v = \sum_{i=1}^{N_k} p_i \phi_k^i$$

be an element of M_k . To apply \mathbf{A}_k to $p = (p_1, \dots, p_{N_k})^t$ we first expand v in the basis $\{\tilde{\phi}_k^i\}$, apply the stiffness matrices for \tilde{M}_k , and finally accumulate $\tilde{A}(v, \phi_k^i)$, $i =$

$1, \dots, N_k$. The application of the stiffness matrix corresponding to the space \widetilde{M}_k with nodal basis $\{\widetilde{\phi}_k^i\}$ is standard. As we shall see, the first and last steps are closely related.

The first step above involves computing the nodal representation of a function v with respect to the basis $\{\widetilde{\phi}_k^i\}$ given the coefficients $\{p_i\}$ appearing in (6.1). Thus, we seek the vector $\widetilde{p} = (\widetilde{p}_1, \dots, \widetilde{p}_{\widetilde{N}_k})^t$ satisfying

$$v = \sum_{j=1}^{\widetilde{N}_k} \widetilde{p}_j \widetilde{\phi}_k^j.$$

Note that $\widetilde{p}_j = p_j$ for $j = 1, \dots, N_k$. Thus, we only need to determine the values of \widetilde{p}_j for the remaining indices. These indices appear in some set $\widetilde{\mathcal{N}}_k^\gamma$ corresponding to one of the interface segments. We define the *transfer matrix* $\mathbf{T}^{k,\gamma}$ by

$$\sum_{j \in \widetilde{\mathcal{N}}_k^\gamma} \mathbf{T}_{ji}^{k,\gamma} \widetilde{\phi}_k^j = \mathcal{E}_{k,\gamma} \widetilde{\phi}_k^i \quad \text{for all } i \in \mathcal{N}_k^\gamma.$$

Then for $j \in \widetilde{\mathcal{N}}_k^\gamma$,

$$\widetilde{p}_j = \sum_{i \in \mathcal{N}_k^\gamma} \mathbf{T}_{ji}^{k,\gamma} p_i.$$

It therefore suffices to have the matrices $\mathbf{T}^{k,\gamma}$ in computations. Note that the entries of $\mathbf{T}^{k,\gamma}$ are obtained by inverting a mass matrix. This inversion can be done once and for all at the beginning of the computations, and the matrices $\mathbf{T}^{k,\gamma}$ can be stored. The storage space required is comparable to that of local stiffness matrices, as each $\mathbf{T}^{k,\gamma}$ is of size $|\mathcal{N}_k^\gamma| \times |\widetilde{\mathcal{N}}_k^\gamma|$.

The last step of accumulating $\widetilde{A}(v, \phi_k^i), i = 1, \dots, N_k$ is also implemented in terms of $\mathbf{T}^{k,\gamma}$. Given the results of the stiffness matrix evaluation on \widetilde{M}_k , i.e., the vector of values $\widetilde{A}(v, \widetilde{\phi}_k^j)$, we need to compute $\widetilde{A}(v, \phi_k^i)$. Clearly, $\widetilde{\phi}_k^i = \phi_k^i$ for nodes which are not on any of the interface segments, so we only need to compute $\widetilde{A}(v, \phi_k^i)$ for nodes such that $i \in \mathcal{N}_k^\gamma$ for some segment. This is given by

$$\widetilde{A}(v, \phi_k^i) = \widetilde{A}(v, \widetilde{\phi}_k^i) + \sum_{\gamma} \sum_{j \in \widetilde{\mathcal{N}}_k^\gamma} \mathbf{T}_{ji}^{k,\gamma} \widetilde{A}(v, \widetilde{\phi}_k^j).$$

The sum on γ above is over the segments with $i \in \mathcal{N}_k^\gamma$.

For convenient notation, let us denote by \mathbf{T}_k the matrix of the linear process that takes $\{p_i : i = 1, \dots, N_k\}$ to $\{\widetilde{p}_i : i = 1, \dots, \widetilde{N}_k\}$. Then the matrix corresponding to $\{\widetilde{A}(v, \widetilde{\phi}_k^i)\} \rightarrow \{\widetilde{A}(v, \phi_k^i)\}$ is the transpose \mathbf{T}_k^t .

We now discuss the implementation of the preconditioner B_k . We use uppercase letters to denote vectors of inner-products and lowercase ones to denote vectors of coefficients. What is needed is a procedure that will compute the coefficients of $B_k g$ in the basis $\{\phi_k^i\}$, given the values $[G_k]_i = (g, \phi_k^i), i = 1, \dots, N_k$. The corresponding matrix will be denoted by \mathbf{B}_k . Clearly, $\mathbf{B}_1 = \mathbf{A}_1^{-1}$. The matrix that takes a vector $\{(w, \phi_k^i)\}$ to coefficients of $R_k w$ with respect to $\{\phi_k^i\}$ will be denoted by \mathbf{R}_k . Finally, let \mathbf{C}_k be the matrix associated with I_k , i.e.,

$$I_k \phi_{k-1}^i = \sum_{j=1}^{N_k} [\mathbf{C}_k]_{ij} \phi_k^j.$$

Assuming \mathbf{B}_{k-1} has been defined, we define $\mathbf{B}_k G_k$ for $G_k \in \mathbb{R}^{N_k}$ by the following four steps:

1. Compute x^l for $l = 1, \dots, m(k)$ by $x^l = x^{l-1} + \mathbf{R}_k(G_k - \mathbf{A}_k x^{l-1})$.
2. Set $y^0 = x^{m(k)} + \mathbf{C}_k^t q_{k-1}$, where $q_{k-1} = \mathbf{B}_{k-1} \mathbf{C}_k(G_k - \mathbf{A}_k x^{m(k)})$.
3. Compute y^l for $l = 1, \dots, m(k)$ by $y^l = y^{l-1} + \mathbf{R}_k(G_k - \mathbf{A}_k y^{l-1})$.
4. Set $\mathbf{B}_k G_k = y^{m(k)}$.

This algorithm is straightforward to implement as a recursive procedure provided we have implementations of \mathbf{R}_k , \mathbf{C}_k , \mathbf{C}_k^t , \mathbf{A}_k , and \mathbf{A}_1^{-1} .

To compute $q_k = \mathbf{C}_k^t q_{k-1}$, we first let $\tilde{q}_{k-1} = \mathbf{T}_{k-1} q_{k-1}$. We then apply the coarse to fine interpolation corresponding to the imbedding $\tilde{M}_{k-1} \subset \tilde{M}_k$. This gives a vector which we denote by \tilde{q}_k . Then q_k is given by the truncated vector $(\tilde{q}_k^1, \dots, \tilde{q}_k^{N_k})^t$.

To compute the action of the transpose, $G_{k-1} = \mathbf{C}_k G_k$, we start by defining \tilde{G}_k to be the vector which extends G_k by $\tilde{G}_k^i = 0$ for $i > N_k$. Next we apply the adjoint of the coarse to fine imbedding $(\tilde{M}_{k-1} \subset \tilde{M}_k)$ to define the vector \tilde{G}_{k-1} . Then $G_{k-1} = \mathbf{T}_{k-1}^t \tilde{G}_{k-1}$.

Since our codes do not assemble matrices, we use the alternative smoother

$$\underline{R}_k g = \Lambda_k^{-1} \sum_{i=1}^{N_k} (g, \phi_k^i) \phi_k^i,$$

where Λ_k is the largest eigenvalue of \mathbf{A}_k . This avoids the computation of the diagonal entry $\tilde{A}(\phi_k^i, \phi_k^i)$. The corresponding matrix operator \mathbf{R}_k is just multiplication by Λ_k^{-1} .

We now show that this operator is a good smoother by showing that it satisfies Condition (C.1). First, (4.2) holds for \underline{R}_k since by Lemma 4.4,

$$\begin{aligned} \frac{\|v\|_{0,\Omega}^2}{\lambda_k} &\leq C_R (R_k v, v) = C_R \sum_{i=1}^{N_k} \frac{(v, \phi_k^i)^2}{\tilde{A}(\phi_k^i, \phi_k^i)} \\ &\leq C \Lambda_k^{-1} \sum_{i=1}^{N_k} (v, \phi_k^i)^2 = C(\underline{R}_k v, v). \end{aligned}$$

Now let v be in M_k and p be as in (6.1). Then

$$\begin{aligned} (\underline{R}_k A_k v, A_k v) &= \Lambda_k^{-1} \sum_{i=1}^{N_k} \tilde{A}(v, \phi_k^i)^2 \\ &= \frac{\mathbf{A}_k p \cdot \mathbf{A}_k p}{\Lambda_k} \leq \mathbf{A}_k p \cdot p = (A_k v, v). \end{aligned}$$

This shows that $I - \underline{R}_k A_k$ is nonnegative and hence Condition (C.1) is satisfied.

7. Numerical results. In this section we give the results of model computations which illustrate that the condition numbers of the preconditioned system remain bounded as the number of levels increase. The program used for computations, takes as input general triangulations generated independently on subdomains, recursively refines these triangulations by breaking each triangle into four similar ones, and solves a mortar finite element problem using the mortar multigrid preconditioner.

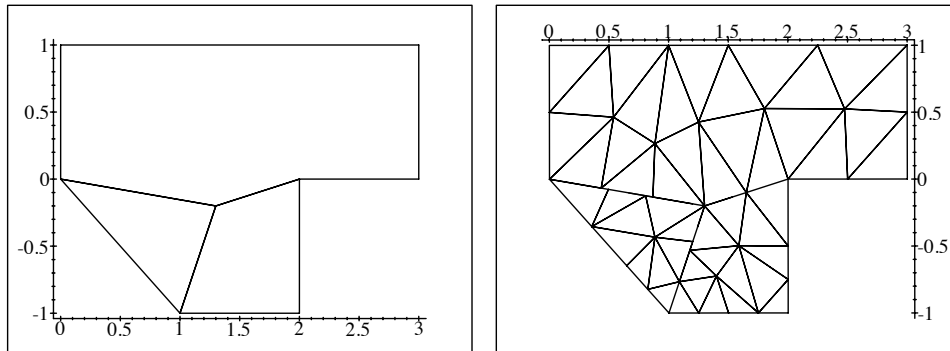


FIG. 7.1. Domain decomposition and initial triangulation for a domain with a re-entrant corner.

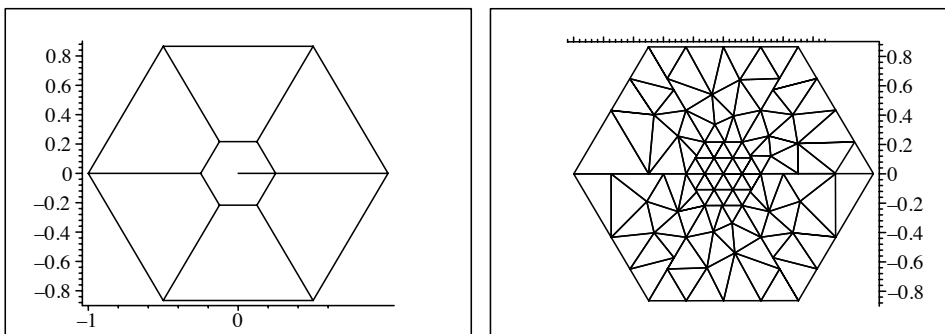


FIG. 7.2. Domain decomposition and initial triangulation for a domain with a cut.

We compute the mortar finite element approximation to a solution of the Poisson equation with zero Dirichlet boundary conditions on two domains. In both cases, the triangulations were done using the mesh generator TRIANGLE [20]. The smoother used was \underline{R}_k defined in the previous section and $m(k) = 2^{J-k}$. Estimates of extreme eigenvalues of the operator $\mathbf{B}_J \mathbf{A}_J$ were given by those of the Lanczos matrix (see [18]). Note that the eigenvalues of $\mathbf{B}_J \mathbf{A}_J$ coincide with those of $B_J A_J$.

Figure 7.1 shows the first domain, its decomposition into three subdomains, and the coarse meshes on subdomains. Note that this domain has a re-entrant corner. By virtue of a theorem in [15], solutions to Poisson equation on this domain satisfy Assumption (A.1) with $\beta < 2/3$.

The domain decomposition and coarse triangulations of the second domain are pictured in Figure 7.2. This hexagonal domain has a cut that extends from the origin to a vertex of the hexagon. The domain is divided into seven subdomains. Assumption (A.1) does not hold for this domain.

Table 7.1 gives the condition numbers of $\mathbf{B}_J \mathbf{A}_J$ for the first domain. As can be seen from the table, the condition numbers remain bounded independently of the number of levels as predicted by the theory. In Table 7.2, we present the results from computations on the second domain. Even though Assumption (A.1) does not hold for this domain, we see that the condition numbers of the preconditioned system remain bounded.

TABLE 7.1
Conditioning of $\mathbf{B}_J \mathbf{A}_J$ for the domain in Figure 7.1.

Level J	Minimum eigen- value of $\mathbf{B}_J \mathbf{A}_J$	Maximum eigen- value of $\mathbf{B}_J \mathbf{A}_J$	Condition number	Degrees of freedom
2	0.59	1.13	1.92	67
3	0.56	1.07	1.90	343
4	0.52	1.09	2.10	1451
5	0.47	1.10	2.34	5971
6	0.45	1.10	2.48	24227
7	0.44	1.10	2.52	97603

TABLE 7.2
Conditioning of $\mathbf{B}_J \mathbf{A}_J$ for the domain in Figure 7.2.

Level J	Minimum eigen- value of $\mathbf{B}_J \mathbf{A}_J$	Maximum eigen- value of $\mathbf{B}_J \mathbf{A}_J$	Condition number	Degrees of freedom
2	0.61	1.32	2.24	221
3	0.54	1.16	2.14	911
4	0.43	1.08	2.73	3731
5	0.38	1.12	3.13	15131
6	0.36	1.12	3.33	60971

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