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# The Influence of Channel Deepening on Tides, River Discharge Effects, and Storm Surge

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# The influence of channel deepening on tides, river discharge effects, and storm surge

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Key Words: Tides, Storm Surge, River flow, Flood Hazard, Compound Flooding, Data Rescue

Key Points:

- Tidal range amplifies after channel deepening in a strongly frictional estuary, with a peak increase near the damping lengthscale.
- Storm surge amplitudes evolve similarly to tides, with a similar spatial pattern and location of maximum change.
- Extreme water levels caused by river discharge have likely decreased due to bathymetric change

## Abstract

We combine archival research, semi-analytical models, and numerical simulations to address the following question: how do changes to channel geometry alter tidal properties and flood dynamics in a hyposynchronous, strongly frictional estuary with a landward decay in tidal amplitudes? Records in the Saint Johns River Estuary since the 1890s show that tidal range has doubled in Jacksonville, Florida. Near the estuary inlet, tidal discharge approximately doubled but tidal amplitudes increased only ~6%. Modeling shows that increased shipping channel depths from 5-6 to ~13m drove the observed changes, with other factors like channel shortening and width reduction producing comparatively minor effects. Tidal amplitude increases are spatially variable, with a maximum change 20-25 km from the estuary inlet; tidal theory suggests that increases in amplitude approximately follow  $x \exp(\mu x)$ , where  $x$  is the distance from the ocean and  $\mu$  is a damping coefficient. Tidal changes are a predictor of altered surge dynamics: Numerical modeling of hurricane Irma under 1898 and 2017 bathymetric conditions confirms that both tidal and storm surge amplitudes are larger today, with a similar spatial pattern. Nonetheless, peak water levels are simulated to be larger under 1898 bathymetry. The cause is likely the record river discharge observed during the storm; as suggested by a subtidal water-level model, channel deepening since 1898 appears to have reduced the average surface slope required to drain both mean river flow and storm flows towards the ocean. Nonetheless, results suggest an increased vulnerability to storms with less river flow, but larger storm surge.

## 37 Plain Text Abstract

38

39 In this study, we evaluate whether channel deepening and other geometric changes have altered  
40 the effects of tides, storm surge, and river flow within the lower Saint Johns River Estuary,  
41 Florida. Using data from archives and old reports, we find that tidal range has more than  
42 doubled in some locations since the late 1800s. Further, the average water level between  
43 Jacksonville, Florida and the coast appears to have decreased, while tidal velocities and  
44 discharge have increased. Numerical and analytical models show that the primary cause is  
45 channel deepening and dredging; other factors, such as shortening the channel, have  
46 comparatively influence. Using the numerical model, we simulated the effects of hurricane Irma  
47 under both modern and historic (1900 era) geometry. Results show that the storm surge from  
48 hurricane Irma was higher today than it would have been a century ago. However, overall water  
49 levels in Jacksonville were simulated to be 0.2 m less today than historically, since the deeper  
50 channel enabled the record amount of rainfall, runoff, and wind-induced currents from the storm  
51 to exit towards the ocean more easily. Hence, anthropogenic development of estuarine  
52 waterways can both decrease the hazard from river-based floods, while increasing the marine-  
53 sourced hazard.

54

## 55 1 Introduction

56 Shipping channels in many estuaries around the world have been deepened by a factor of two or  
57 more since the mid-19<sup>th</sup> century, with deep-draft ships requiring increasingly wide and deep  
58 shipping channels (e.g., Winterwerp et al., 2013, Talke & Jay, 2020). At the same time,  
59 channelization, reclamation and diking has often reduced connectivity to wetlands and reduced  
60 estuary width. Consequences include increased salinity intrusion (e.g., Ralston & Geyer, 2019),  
61 altered tidal velocities (e.g., Pareja-Roman et al., 2020) and an upstream movement of the  
62 estuary turbidity maximum (see review by Burchard et al., 2018, and references therein).  
63 Reduced frictional resistance in a deeper channel leads to reduced damping of long-wave energy.  
64 Depth changes also alter resonance effects and can lead to either amplification or attenuation of  
65 tidal amplitudes (Talke & Jay, 2020), particularly if a total reflection occurs at the head of an  
66 estuary. Convergence and width changes also influence tidal amplitudes (Jay, 1991; Friedrichs  
67 and Aubrey, 1994).

68 The combination of frictional and resonance effects have sometimes resulted in a doubling (or  
69 more) of tidal range in up-river locations (Di Lorenzo et al., 1993; Winterwerp & Wang, 2013;  
70 Talke & Jay, 2020). Since storm surge is a shallow-water wave with a similar amplitude and time  
71 scale as a tide wave, the same anthropogenic alterations can produce similar increases in storm  
72 surge magnitudes (Famalkhalili & Talke, 2016; Ralston et al., 2019). Nonetheless, decreased  
73 frictional effects can lead to a decrease in mean (tidally averaged) water levels for a given river  
74 flow, due to a reduced subtidal slope in the water surface (Jensen et al., 2003; Jay et al., 2011;  
75 Ralston et al., 2019). Because the same process (channel deepening) can amplify long-wave  
76 heights but decrease the tidally averaged water level, it remains unclear whether channel

77 deepening will produce *higher* or *lower* total water levels for any given storm event at a given  
78 location (combined tides, surge, local wind setup, precipitation and river flow).

79 In this contribution we study an estuary known to be sensitive to channel dredging. Numerical  
80 simulations of the lower Saint Johns River Estuary (SJRE) suggest that future sea-level rise and  
81 planned channel deepening to 14.3 m will likely increase tidal range on the order of 0-0.1 m  
82 within the channel (Bilskie, 2013; Hagen et al., 2013), and increase the 50-100 year storm tide by  
83 0-0.2 m (USACE, 2014). These changes, along with other environmental effects such as  
84 increased salinity intrusion (Bellino & Spechler, 2013; Mulamba et al., 2019), show that the  
85 region is sensitive to anthropogenic modification. Given that historical shipping channel depths  
86 have increased from perhaps 3 m pre-1870 to ~12.2 m today, the SJRE is a good test bed for  
87 examining mechanisms of change and likely results of deepening.

88 The use of the Saint Johns River estuary as a case study is further motivated by a historical  
89 conundrum. On October 2, 1898, a category 4 hurricane made landfall just north of the  
90 Florida/Georgia border, severely flooding the town of Fernandina (Monthly Weather Review,  
91 1898) and producing a water level of 2.6 m (relative to the NAVD-88 datum) at the mouth of the  
92 Saint Johns River, located 35km south (Figure 1; Sandrik and Jarvinen, 1999). Nonetheless,  
93 damage in the city of Jacksonville, located ~40km inland along the Saint Johns River, was minor  
94 (Monthly Weather Review, 1898). By contrast, storm surge from hurricane Irma on Sept. 11,  
95 2017 caused record flooding in Jacksonville and nearby regions (e.g., Monroe & Hong, 2018),  
96 even though the water levels of 1.7 m measured at the estuary inlet (relative to NAVD-88 datum)  
97 were significantly lower than in 1898. Though many factors influence flooding, including the  
98 built environment and the meteorological characteristics of each particular storm (rainfall, storm  
99 track, wind field, and pressure), we focus here on the water level effects of bathymetric change.  
100 Specifically, we ask three related questions:

- 101 • Did channel deepening, streamlining, and other anthropogenic changes to the Saint Johns  
102 River reduce the natural protection against storm surge that the shallower channel of  
103 1898 provided, increasing the vulnerability of inland regions to marine-sourced flooding?
- 104 • Did these same changes facilitate drainage of precipitation run-off to the ocean, reducing  
105 the hazard of the river flood wave?
- 106 • What is the net effect of these landscape changes on water levels, during both typical and  
107 storm conditions?

108 Our case-study approach provides insights into how tides, tidal discharge, mean water levels, and  
109 storm surge in similar hyposynchronous estuaries—highly frictional and marked by tides that  
110 strongly decay in the landward direction—might react to anthropogenic modifications such as  
111 channel deepening. We employ a combination of archival data rescue, semi-analytical modeling,  
112 and numerical modeling to obtain new insights into the long-term trajectory of change, and their  
113 causes. Archival tidal records from as early as the 1890s are digitized and used to quantify  
114 spatially variable changes to tidal range and estimate river slope. A semi-analytical model is used  
115 to explore how depth and other geometric variations influence both tidal and subtidal properties.  
116 Finally, we use numerical simulations based on 1898 and 2014 bathymetry to explore how storm  
117 surge and peak water levels during hurricane Irma (2017) were affected by bathymetric change.

## 118 2 Setting and Methods

119

### 120 2.1 Setting and bathymetric change

121 The Saint Johns River, Florida, is a microtidal estuary with primarily semidiurnal tides. Tidal  
122 range decreases from ~1.5 m at the ocean inlet to approximately 0.55 m in the city of  
123 Jacksonville, located 40km upstream. The estuary is heavily channelized and diked over its first  
124 40km. A large, shallow bay extends southward from Jacksonville for 85 km, with a typical width  
125 between 2-5km and a controlling channel depth of 4.5 to 6 m (Figure 1; NOAA chart 11492).  
126 Tides propagate southwards from Jacksonville along the tidal river until finally dissipating more  
127 than 100km upstream (Henrie & Valle-Levinson, 2014). The average river discharge from the  
128 nearly 23,000 km<sup>2</sup> watershed for the years 1971-2017 was ~200 m<sup>3</sup>/s (see also Bellino &  
129 Spechler, 2013). A record tidally filtered flow of just over 4000 m<sup>3</sup>/s was estimated by the  
130 United States Geological Survey (USGS) on Sept. 12<sup>th</sup>, 2017 (USGS gauge 02246500),  
131 approximately one day after the peak storm surge from hurricane Irma.

132 Over the past 150 years, the shipping channel in the Saint Johns River has been deepened from a  
133 controlling depth of 6-10 feet (1.8-3 m) to a depth that varies between 41 to 50 feet (12.5-15.2 m;  
134 NOAA Chart 11491-02-2015) relative to Mean-Lower-Low Water (MLLW). Modest dredging  
135 that totaled ~ 75-80,000 cubic meters was begun in the 1870s to develop and maintain a 10 ft.  
136 (3m) deep, 80 ft. (24.4 m) wide shipping channel (Kingman et al., 1915). Construction of an  
137 entrance jetty to scour the mouth began in the 1880s, and the jetties currently extend ~2 km into  
138 the ocean. Such channelization and dredging efforts increased the controlling depth to 3.7-4.6 m  
139 (12-15) feet in the early 1890s (USACE, 1893). By 1900, a shipping channel of 18ft. (5.5 m) was  
140 dredged nearly to Jacksonville (e.g., NOAA chart 454A-12-1899), and diking of wetlands had  
141 begun. The scale of dredging increased in the early 20<sup>th</sup> century, and by 1915 the channel to  
142 Jacksonville was reported to be 7.9-9.1 m (26-30 ft.; NOAA chart 577-00-1917). In 1952, the  
143 shipping channel was shortened by 3 km by cutting through a wetland located just downstream  
144 of Dames Point (See Figure 1). An additional kilometer of length was removed near the  
145 Longbranch neighborhood of Jacksonville at River kilometer (Rkm) 30 (see Figure 1b with 1c,  
146 or Supplemental Information.). By 1959, channel depths varied from 9.1 m (30 ft., in  
147 Jacksonville) to 12.2 m (40 ft., at entrance) relative to MLLW, with a dredged width between  
148 122 to 366 m (400 to 1200 ft.). A history of changes is available in Rawls (1952). Plans are  
149 currently being implemented to further dredge to 14.3 m (47 ft.; USACE, 2014).

150 To enable numerical modeling of the 1890s era (see Section 2.5), we digitized available  
151 bathymetry from 1898 from the coastal ocean to just upstream of Jacksonville (NOAA chart  
152 454A-12-1899). For modern bathymetry, we obtained a digital elevation model from 2014 from  
153 NOAA/NOS (National Oceanic and Atmospheric Administration/National Ocean Survey) (see  
154 Figure 1). Further information about the location and depth of tributary channels was taken from  
155 additional charts available at <https://historicalcharts.noaa.gov> . Because both historical and  
156 modern bathymetric surveys typically only include data below Mean Low Water, wetland  
157 bathymetry from a Lidar survey with 5 m resolution was obtained from NOAA  
158 (<https://coast.noaa.gov/digitalcoast/data>). The areal extent of wetlands near the coast appears to  
159 be similar in both the historical and modern maps, except near the industrial corridor around the  
160 shipping channel. Therefore, due to a lack of historical wetland bathymetry, we use the modern  
161 Lidar data as a proxy for the historical floodplain.

162 Based on an analysis of the digital elevation models, the flow-carrying width has been narrowed,  
163 with the median width in the modern system about 40% less than in the 1890s (Figure 2). At the  
164 same time, the cross-sectionally averaged depth has slightly more than doubled (Figure 3). The  
165 measured width change is due to land reclamation and diking; over the same time period, the  
166 width of the shipping channel increased from xxx to yyy m. The difference between the total  
167 width and the shipping channel width stems from our definition of the flow-carrying width  
168 (Figure 2), which we measure between the Mean Low Water lines at the side of the channel. The  
169 average depth was calculated by dividing the cross-sectional area by the width (Figure 3).

170

## 171 2.2 Data

172 We use a combination of historical and modern records to elucidate the effects of bathymetric  
173 changes on tidal properties. A continuous hourly record of water level is available from Mayport  
174 since 1928, from a composite of NOAA gauges 8720218 and 8720220. Modern NOAA data for  
175 Dames Point (station number 8720219; 1998-present, intermittent), the Longbranch  
176 neighborhood of Jacksonville (station 8720242; 1998-2003, intermittent) and Jacksonville  
177 (station 8720226; 1997-present, intermittent) are also used in our analysis (see Figure 1 for  
178 locations). The USGS has also measured water levels and discharge near the Acosta Bridge in  
179 Jacksonville since 1945, but only records since 1970 are available (station 02246500). Though  
180 the USGS and NOAA gauges in Jacksonville are less than 1km apart, they measured peak water  
181 levels that were approximately 0.15 m different (relative to NAVD-88 datum) during Hurricane  
182 Irma. The reason is unknown, but could include gauge error or substantial local variability in  
183 water levels.

184 Archives and old reports yielded substantial information about historical tidal conditions (see  
185 Supplemental Information Section S.2). Synopses of tidal measurements from 1879, 1889, and  
186 1892 were found in the annual reports of the United States Army Corps of Engineers (USACE  
187 1879, 1892; Gieseler, 1893). Additional extracts of measurements taken between the 1850s and  
188 early 1900s were found in summary sheets of the United States Coast and Geodetic Survey  
189 (USCG&S, Record Group 23, United States National Archives). Monthly mean tide level and  
190 mean tidal range for Mayport, Florida for 1895-1897 are available from NOAA (station  
191 8720220). Historical measurements of tidal range are converted to an  $M_2$  estimate by dividing by  
192 2.07, the ratio observed in modern measurements. An estimate of mean water level for  
193 Jacksonville for the year 1892 relative to the North American Vertical Datum of 1988 (NAVD-  
194 88) was found using the tabulated height of mean high and low water relative to local  
195 benchmarks (see Supplemental Information Section S.2.4).

196 Hourly and high/low tidal records were also found, recovered, digitized, and quality assured (see  
197 also Talke & Jay, 2017), as summarized in Supplemental Table S.1. Three years of hourly data  
198 from Mayport (1895-1897) were recovered from the United States National Archives and  
199 digitized. Tidal records from 1928-1935 and 1953-1968 from Longbranch (Rkm 30) were found  
200 in the EV2 database from the National Centers for Environmental Information  
201 (<https://www.ncdc.noaa.gov/EdadsV2/>); selected high/low data and summary statistics such as  
202 mean tidal range and mean sea level were digitized. The datum for the 1929-1935 and 1953-  
203 1968 series was tied to the NAVD-88 datum through an extant benchmark (see supplement  
204 S.2.3). A short NOAA record from Acosta Bridge for the year 1959 was also digitized.



205 Historical estimates of cross-sectionally-averaged discharge at 6 locations are available between  
206 Mayport and Rkm 30 from Gieseler (1893), based on 9 days of measurements at Mayport  
207 between August 12 and August 23, 1892. Estimates upstream of Mayport were approximated by  
208 calculating tidal prism from tidal measurements. We adjusted the averaged flood/ebb discharges  
209 to an  $M_2$  amplitude by assuming an equivalent sinusoidal discharge. Although these  
210 measurements must be considered approximate, they show good agreement with modeling (see  
211 Results) and therefore help ground-truth results. A modern estimate of the  $M_2$  discharge  
212 amplitude was obtained through harmonic analysis (Pawlowicz et al., 2002) of USGS discharge  
213 data at Jacksonville over the period of record.

214

### 215 2.3: Semi-analytical tidal model

216 To help interpret water level changes evident in archival records (see results), we develop an  
217 idealized analytical tidal model (this section) and subtidal model (section 2.4). The width, depth  
218 and planform of both models are presented in Figures 2-4. The modern and historical system  
219 between the ocean and Jacksonville are approximated as two constant width and depth channels  
220 (see dashed lines in Figure 2 & 3), and reflect the observation that there is no clear depth or  
221 width convergence within the system. Further, the use of constant width and depth facilitates  
222 comparison with the subtidal water level model discussed below (Section 2.4). In the historical  
223 configuration, a 4 m deep channel transitions to a 6 m deep channel upstream of Rkm 28. Width  
224 is held constant at 1500m. In the modern configuration, a 10m deep channel seaward of Rkm 20  
225 transitions to 8 m deep between Rkm 20-48. The width is somewhat wider near Jacksonville  
226 than the channelized coastal section (1100m vs. 800m). Upstream (south) of Jacksonville, a long,  
227 shallow region (4 m depth, 4km width, 85km long) is modeled. A gradual transition to this wide  
228 channel is applied. A long, 100+ km narrow section that resembles the observed tidal river is  
229 included upstream of the shallow bay. The shallow upstream regions are required to allow the  
230 tide wave to decay towards zero, and to reproduce the observed tidal discharge, tidal prism, and  
231 tidal phases (see e.g. Wang et al. 2019). Due to channel streamlining, the modern planform is 4  
232 km shorter than the historical planform (see section 2.1).

233 A linearized, semi-analytical tide model is employed to gain insights into the reasons for tidal  
234 change. Schematized analytical models have often been used to explore how depth and other  
235 parameters affect tidal amplitudes ( e.g., Jay 1991; Friedrichs & Aubrey, 1994; Winterwerp et  
236 al., 2013). Our analysis follows Dronkers (1964), makes the shallow-water approximation (wave  
237 length long compared to depth), and assumes that the tide wave amplitude  $\eta$  is small relative to  
238 depth. Based on observations (see results), depth-averaged velocities are dominated by the  $M_2$   
239 tide, to leading order; subtidal velocities are more than an order of magnitude smaller during  
240 typical conditions, and are neglected here. We set our coordinate system at the ocean boundary,  
241 and let  $x$  be positive into the estuary. Assuming constant width  $b$  and depth  $h$ , the depth and  
242 width integrated mass and momentum balance within a section are

$$243 \frac{\partial Q_t}{\partial x} + b \frac{\partial \eta}{\partial t} = 0, \quad (1)$$

244

245  $\frac{\partial Q_t}{\partial t} + rQ_t + gbh \frac{\partial \eta}{\partial x} = 0,$  (2)

246

247 where  $Q_t$  is the tidal discharge,  $g$  is gravity, and  $r$  is the linearized frictional resistance. Under the  
 248 assumption that tidal discharge is much larger than river discharge ( $Q_t \gg Q_r$ ), the linearized  
 249 friction coefficient can be approximated (using the first term of a Fourier or Chebyshev  
 250 expansion of  $Q_t|Q_t|$ ; Dronkers, 1964) as

251  $r = \frac{8}{3\pi} \frac{C_d Q_T}{bh^2},$  (3)

252 where  $Q_T$  is the tidal discharge amplitude and  $C_d$  is the drag coefficient. Following the solution  
 253 procedure described in Dronkers (1964), a solution of the following form can be derived:

254  $\eta(x, t) = \text{Re} \left[ \left( \underbrace{A_0 e^{kx}}_{\text{Reflected wave}} + \underbrace{B_0 e^{-kx}}_{\text{Incident Wave}} \right) e^{i\omega t} \right],$  (4)

255 where  $A_0$  and  $B_0$  are constants for the reflected and incident wave components. The frequency  $\omega$   
 256 is related to the tide period  $T$  by  $\omega = \frac{2\pi}{T}$ , and  $k$  is a complex number described by

257  $k = \frac{\omega}{\sqrt{gh}} \left( -1 + \frac{ir}{\omega} \right)^{1/2}.$  (5)

258 An equation for tidal discharge  $Q_t(x, t)$  then follows from the continuity equation (Equation 1).  
 259 The solution for  $\eta(x, t)$  and  $Q_t(x, t)$  is found by applying boundary conditions. At the ocean  
 260 boundary, we apply a sinusoidal wave at the  $M_2$  frequency ( $T = 12.42$  hours) with an amplitude  
 261 of  $\eta_o = 0.7$  m. At the upstream boundary, a no-flux condition is applied. Following Dronkers  
 262 (1964), we further subdivide the model into  $N$  segments, each of 4 km length. For the internal  
 263 boundaries between segments, the tidal discharge  $Q_t(x, t)$  and water level  $\eta(x, t)$  at the upstream  
 264 boundary of each segment is matched to the downstream boundary condition of the next  
 265 segment. This produces a system of  $2N$  equations which is solved through matrix inversion. The  
 266 tidal amplitude  $\eta$  and discharge  $Q_T$  is solved iteratively. First, a solution is found using an initial  
 267 estimates for  $Q_T$ . The friction term (Equation 3) is re-calculated using updated estimates of  $Q_T$ ,  
 268 and the equations re-solved. The solution is iterated until it changes by less than 0.1% between  
 269 successive approximations.

270 The model was calibrated by changing the value of the drag coefficient  $C_d$  and comparing the  
 271 solution with measured tidal amplitudes and discharge (see Results). To avoid coding errors, we  
 272 also checked the model against the analytical solution of a constant width and depth  
 273 configuration (Dronkers, 1964). Through calibration, the optimal drag coefficient for the  
 274 historical and modern configuration was  $C_d = 0.007$  and  $C_d = 0.005$ , respectively, within the range  
 275 of 0.001 to 0.01 typically found for analytical models (Friedrichs & Madsen, 1992). The  
 276 equivalent Manning's  $n$  roughness coefficient is 0.031-0.033 s/m<sup>1/3</sup> (modern configuration) and  
 277 0.033-0.035 s/m<sup>1/3</sup> (historical configuration), using the conversion formula  $n = R^{1/6} \left( \frac{C_d}{g} \right)^{1/2},$



278 where  $R$  is the hydraulic radius (Area divided by wetted perimeter) and is approximately equal to  
 279  $h$  in a wide channel. A root mean square error (RMSE) of 0.035 and 0.044 m was found between  
 280 measurements and the modern and historical configurations, based on 5 and 8 measurements,  
 281 respectively (see Results and Table 1). Tidal discharge amplitudes for both configurations  
 282 agreed with measurements to within 10%, and the relative phase between discharge (velocity)  
 283 and water level agreed well, to within 10 degrees (Table 1). The analytical model estimated a  
 284 phase progression of 44 and 30 degrees between Mayport and Jacksonville for the historical and  
 285 modern configuration; the equivalent based on available measurements was 40 and 49 degrees.

286

## 287 2.4: Subtidal water level from river discharge

288

289 We next develop an analytical model for how the tidally-averaged (subtidal) water surface is  
 290 influenced by geometry changes. The same geometry as in the tidal model (section 2.3) is  
 291 considered. As shown by Godin (1999), the effective subtidal friction is set by both river flow,  
 292 tidal forcing, and non-linear interaction between both (see also Kukulka & Jay, 2003a;  
 293 Buschman et al.; 2009 ). Because average river discharge (200 m<sup>3</sup>/s) is small compared to the  
 294 typical M<sub>2</sub> tidal discharge (~4200 m<sup>3</sup>/s), we follow Godin (1999) and Buschman et al. (2009) and  
 295 examine the parameter space in which tidal currents outweigh river flow currents. We also  
 296 assume that bed slope effects and the effect of surface slope on water depth are negligible.  
 297 Results show that subtidal water level variations in the estuary region are small relative to mean  
 298 depth under normal conditions, justifying this assumption (see also Henrie & Valle-Levinson,  
 299 2014). A more thorough treatment of bed and river slope effects, particular in tidal rivers, is  
 300 presented in Kästner et al. (2019).

301 For a constant width segment of an estuary, the tidally and sectionally averaged momentum  
 302 balance becomes a balance between the barotropic pressure gradient and tidally-averaged bed  
 303 stress (e.g., Kukulka & Jay, 2003b, Buschmann et al., 2009):

304

$$305 \quad gh \frac{\partial \langle z_r \rangle}{\partial x} = \frac{-\langle T_r \rangle}{\rho}, \quad (6)$$

306

307 where the bed stress is  $T_r$ , the density of water is  $\rho$ , the tidally-averaged surface slope relative to  
 308 a fixed datum is  $\frac{\partial \langle z_r \rangle}{\partial x}$ , and the brackets denote a tidal average. For simplicity, we neglect small  
 309 tributaries and the subtidal discharge caused by the correlation between vertical and horizontal  
 310 tidal velocities (Stokes transport). Equation (6) is simplified by using the definition for bed  
 311 stress,  $T_r = \rho C_d \langle |U|U \rangle$ , where  $U$  is the velocity,  $C_d$  is the drag coefficient, angle brackets denote  
 312 an average over the tide and the absolute value preserves the directionality of stress within the  
 313 brackets. The velocity  $U$  consists of tidal fluctuations and river flow, i.e.,  $U = U_T \cos(\omega t) + U_r$ ,  
 314 where  $U_r$  is negative because discharge moves in the minus  $x$  direction. Following Dronkers  
 315 (1964), we apply a Chebyshev polynomial expansion on the velocity term  $|U|U$ . We then

316 tidally average the expansion term and retain only components that are significant (see  
 317 supplemental information, Kukulka & Jay (2003b) and Buschman et al. (2009). The tidally  
 318 averaged bed stress is then approximated as:

319

$$320 \frac{\langle T_r \rangle}{\rho} = \frac{-C_d}{\pi} \left( p_1 U_o U_R + \frac{3}{2} p_3 U_R U_o \left( \frac{U_T}{U_o} \right)^2 \right), \quad (7)$$

321 where  $p_1 = 16/15$  and  $p_3 = 32/15$  are expansion coefficients, and  $U_o$  is a velocity scale. We  
 322 have defined the positive river velocity scale  $U_R = -U_r$ ; the minus sign in Equation 7 accounts  
 323 for the fact that river discharge moves in the minus  $x$  direction. We follow Buschman et al.  
 324 (2009) and set the velocity scale  $U_o$  to the absolute value of the maximum velocity. Applying  
 325 the simplifications discussed above, we find that the differential equation for  $\frac{\partial \langle z_r \rangle}{\partial x}$  can be  
 326 approximated as:

327

$$328 \frac{\partial \langle z_r \rangle}{\partial x} = 1.36 \frac{C_d Q_R U_T}{g b h^2}, \quad (8)$$

329

330 where  $Q_R = b h U_R$  is the river discharge. More generally, since  $p_1$  and  $p_3$  change slightly as the  
 331 ratio of river to tidal discharge varies (e.g., due to spring-neap cycle), we state that, to first order,  
 332  $\frac{\partial \langle z_r \rangle}{\partial x} \sim \frac{C_d Q_R U_T}{g b h^2}$ . A similar, more complex analysis of the subtidal slope that includes additional  
 333 tidal bands is found in Buschman et al., 2009.

334 The analytical development in Equations 6-8 is applied to the geometry of our model (Figures 2-  
 335 4) by requiring water level at the boundary of each constant width/depth segment (of 4 km  
 336 length) to match the next. We integrate Equation 8 under the assumption that the bed slope is  
 337 negligible and that river discharge, tidal velocity, width, and depth are approximately constant  
 338 over the 4km segment under consideration. This yields

339

$$340 \langle z_r(x) \rangle = \alpha x + z_{ri}, \quad (9)$$

341

342 where  $\alpha = 1.36 \frac{C_d Q_R U_T}{g b h^2}$  is the subtidal river slope and the constant of integration  $z_{ri}$  is the mean  
 343 water level at the downstream boundary of segment  $i$ . At the ocean ( $x=0$ ), sea-level is used as a  
 344 boundary condition. For consistency, the subtidal model uses the drag coefficients  $C_d$  that were  
 345 calibrated from the historical and modern analytical tidal model ( $C_d = 0.007$  and  $C_d = 0.005$ )

346

## 347 2.5 Delft-3D Numerical model

348 To assess the storm tide produced by hurricane Irma under historic and modern channel  
349 conditions, we run simulations using the Delft3D numerical model (Deltares, 2014). Such a  
350 model is better suited for modeling unsteady, energetic storm conditions than the simplified  
351 analytical models described above, and better represents system depth. Two configurations were  
352 developed, one using a grid based on 1898 bathymetry (see section 2.1), and another based on  
353 modern bathymetry from 2014 (Figure 1). The domain stretches from the coastal ocean to a  
354 location 100 km upstream of Jacksonville (see Figure 1a), and is divided into two domains  
355 (labeled *A* and *B* in Figure 1). The river upstream of Jacksonville (Domain B) is approximated as  
356 a long, wide and shallow bay to allow the observed damping of the tidal wave. The width is 4000  
357 m and the depth is approximated to be 4 m deep. Bathymetry is not adjusted for sea-level rise.  
358 The model contained 545,500 grid cells, with the majority (~97%) in Domain A between the  
359 ocean boundary and Jacksonville. A grid resolution of 30 m was applied within the shipping  
360 channel.

361 For calibration, the model is run for 40 d using average river discharge ( $200 \text{ m}^3/\text{s}$ ). Tidal forcing  
362 at the boundary is obtained from the NOAA gauge at Mayport and is scaled by a factor of 1.06 to  
363 account for the attenuation of tides through the jetties. A different Manning's friction coefficient  
364 is applied to vegetated and unvegetated parts of the domain, following observations in similar  
365 (local) modeling efforts (Bacopoulos et al., 2012). Following Bacopoulos et al. (2009, 2017a),  
366 we run the model in depth-averaged mode, since the estuary is likely to be well-mixed during  
367 highly energetic storm conditions. In other estuaries, neglecting density variations produces a  
368 small (generally <10%) underestimation in storm tide amplitudes (e.g., Orton et al., 2012); here,  
369 we assess the validity of our approach through comparison with measurements.

370 We calibrate the model to reproduce the observed tidal statistics between the estuary inlet and  
371 Jacksonville (see Results, Table 1). The optimal Manning's  $n$  coefficient for the historical  
372 channel and wetland was  $n = 0.025 \text{ s/m}^{1/3}$  and  $n = 0.05 \text{ s/m}^{1/3}$ , and for the modern configuration  
373 was  $n = 0.02 \text{ s/m}^{1/3}$  and  $n = 0.04 \text{ s/m}^{1/3}$ . Within Domain B, a constant  $n = 0.025 \text{ s/m}^{1/3}$  is used  
374 for both configurations. Simulations agree well, overall, with available measurements.  
375 Simulated tidal discharge agrees to within 2 and 15% with modern and historical measurements  
376 (Table 1), likely within the uncertainty of measurements. Modern simulations and measurements  
377 both depict a progressive wave which takes about 1.5 hours to travel from Mayport and  
378 Jacksonville, with minor differences in phase progression (8 degrees) and relative phase of water  
379 level and tidal discharge (<12 degrees). A somewhat larger difference is observed in the  
380 simulated historical progression of the  $M_2$  wave (24 degrees), likely in large part because of  
381 uncertainty in the empirical estimate (which was estimated from the mean tabulated travel time)

382 Tidal amplitudes are well calibrated in both simulations (Table 1). The root-mean-square error  
383 (RMSE) in the  $M_2$  constituent was 0.025 m (8 measurements) and 0.008 m (5 measurements) for  
384 the historical and modern configurations, respectively. The slightly larger Manning coefficient  
385 historically may reflect larger sub-gridscale roughness (e.g., sand dunes and other bathymetric  
386 variation), or may account for uncertainty in the historical bathymetric measurements.  
387 Conversely, salinity stratification within the modern system (e.g., Bellino & Spechler, 2013;  
388 Bacopoulos et al., 2017b) may also reduce the effective, depth-averaged frictional effect, as has  
389 also been observed at other locations (Giese & Jay, 1989). The RMSE in the historical  
390 configuration only increases to 0.036 m from 0.025 m when the Manning's  $n$  is decreased from

391  $n = 0.025 \text{ s/m}^{1/3}$  to  $n = 0.02 \text{ s/m}^{1/3}$  Therefore, changes to the friction coefficient exert only a  
392 minor influence on tidal results.

393 To simulate hurricane Irma effects on water level, we apply the known water-level variations  
394 during the storm at the ocean boundary (approximately 10km from the estuary inlet; see Figure  
395 1), using data from near the estuary entrance (tide gauge at Mayport). Data are scaled up by 6%  
396 to account for the decay in tides between the boundary and the tide gauge. A similar ‘storm surge  
397 hydrograph’ approach is used in other studies (e.g., Xu & Huang, 2014). Fluvial discharge  
398 effects during Irma are modeled using two approaches. First, we run the storm surge model by  
399 applying a constant discharge of 0 to 7,000  $\text{m}^3/\text{s}$  at the upstream boundary, in increments of 1000  
400 or 2000  $\text{m}^3/\text{s}$ . This enables us to estimate the sensitivity of peak water level to discharge.  
401 Additionally, we also model the discharge measured at USGS station #02246500 in Jacksonville  
402 using a “virtual” boundary condition (Deltares, 2014) This virtual boundary condition forces the  
403 model to reproduce the total discharge measured at Jacksonville (tides + surge+ discharge) by  
404 applying either a source or sink discharge at the gauge location (as needed). We found this  
405 approach greatly improves comparison of the model to measured water levels, compared to using  
406 the USGS ‘tidally filtered’ discharge product at the model boundary. The virtual boundary  
407 approach is needed because the measured discharge includes storm surge currents and the effects  
408 of local winds, which cannot easily be separated from local run-off. Accounting for such factors  
409 requires hydro-meteorological modeling (e.g., as done in Bacopoulos et al. 2017a), and is beyond  
410 the scope of the current effort. Since the virtual boundary approach likely introduces some  
411 uncertainty into the historical discharge forcing, we compare results with the constant discharge  
412 simulations; as shown in the results, the different approaches yield broadly consistent results.

413 The individual effects of tides, storm surge, and local discharge effects on water levels are  
414 decomposed by running 40d “tide-only”, “tide + surge”, and “tide + surge + discharge” model  
415 runs. The differences between these model runs are used to infer the difference between  
416 individual contributions to the total water level. For example, surge effects are estimated by  
417 subtracting “tide-only” results from the “tide + surge” results, and discharge effects are estimated  
418 by subtracting the “tide + surge” model results from the “tide + surge + discharge” model  
419 results. This approach, though commonly used (e.g., Shen et al., 2006), does not account for the  
420 modification of the tidal phase speed by the surge, or non-linear frictional interaction (see e.g.,  
421 Horsburgh and Wilson, 2007; Valle-Levinson et al., 2013; Familkhalili et al., 2020). Hence,  
422 some tidal energy is likely aliased into our estimated surge signal, and some tidal and surge  
423 effects are aliased into our local discharge estimate. For this reason, it is important to check that  
424 numerical simulation results are consistent with available empirical records and that trends are  
425 consistent with analytical and numerical model results obtained during low-discharge conditions.

426

### 427 3 Results

428 We next use our data, analytical modeling, and numerical simulations to explore how tidal  
429 dynamics and the river slope in the Saint Johns River Estuary have shifted during typical  
430 discharge conditions (section 3.1 and 3.2). After discussing measurements during hurricane Irma  
431 (section 3.3), we use numerical simulations to explore how anthropogenic modifications may  
432 have altered water levels during hurricane Irma (section 3.4). Reasons for changes are explored  
433 in the Discussion.

434

### 435 3.1 Tide changes

436 Water level observations depict a continually evolving tidal range over the past century (Figure 5  
437 and 6). Trends near the estuary mouth are small; at Mayport (Rkm 5.5), the tidal range has  
438 increased at a rate of  $\sim 0.33$  mm/yr. since 1892, for a total increase of 0.04 m (Figure 5a), and the  
439  $M_2$  amplitude (Figure 6a) has increased from 0.63 to 0.67 m. By contrast, tidal range from  
440 Dames Point (Rkm 17.3) to the end of the maintained shipping channel ( $\sim$ Rkm 38) has more  
441 than doubled (Figure 7). At Longbranch (Rkm 31; open circles), tidal range increased from 0.33  
442 to 0.77 m since 1900, at an average rate of 5.2 mm/yr.; in downtown Jacksonville (orange stars),  
443 tidal range increased from about 0.29 m to 0.55 m over the same period, an increase of  $\sim 90\%$   
444 (Figure 5b,7). The divergence in trends between stations near the coast and inland stations  
445 points to a local cause, rather than far field changes in the Atlantic Ocean.

446 The observed tidal amplitude changes are reproduced by both the analytical and numerical  
447 models (Figure 6a, Table 1). The maximal tide change occurs within the mid-estuary, roughly  
448 between Rkm 20-30 (Figure 6a, Figure 7). The magnitude of increase becomes less pronounced  
449 further upstream (Figure 6a), even though the percentage increase is still large (Figure 7). Both  
450 modeling and measurements suggest that the tidal discharge amplitude approximately doubled  
451 since the 1890s. The tides have retained their progressive wave characteristic, with tidal flow  
452 nearly in phase with water level (Table 1). The phase progression of the tide wave is about the  
453 same; Gieseler (1893) reports that the tide wave took slightly less than 1.5 hours to propagate  
454 between Mayport and Jacksonville. The time today is  $\sim 1.7$  hours.

455 The reasons for tidal changes are explored by applying sensitivity tests to the analytical model  
456 (Figure 8). Keeping all other parameters equal, we find that increasing the depth produces the  
457 largest increase in tidal statistics. The change is spatially variable, with a peak value of nearly  
458 0.28 m found between Rkm 25 to Rkm 30. Both the decrease in the drag coefficient and channel  
459 length changes produce minor (less than +0.05 m) changes in tidal range, consistent with  
460 numerical modeling results (Figure 6). Shortening effects are cumulative and most prominent  
461 near Jacksonville. Decreasing channel width reduces tidal amplitudes, all other parameters held  
462 equal. The maximum decrease of slightly more than 0.05 m is found around Rkm 20, and may  
463 be caused by the increase in tidal currents (and therefore friction factor  $r$ ) that occurs when width is  
464 decreased. Depth and the drag coefficient both influence tidal amplitudes by altering the friction  
465 factor (Equation 3); since the percentage depth increase is much bigger than the drag coefficient  
466 decrease, its observed influence is larger (Figure 8; see also discussion).

467

### 468 3.2 Mean water level changes

469

470 Archival records show that sea-level is rising both at the estuary inlet and in Jacksonville (Figure  
471 9). At Mayport, sea level has increased at an average rate of 2.5 mm/yr. (Figure 9a) since 1895,  
472 slightly larger than the 2.1 mm/yr. registered 35 km northwards in Fernandina (see Figure 1 for  
473 location). During the 20<sup>th</sup> century, a smaller sea-level rise rate is observed in Jacksonville,

474 compared to Mayport. The difference between the two locations between 1929 and 1995 was  
475  $\sim 1.5 \pm 0.3$  mm/yr. (Figure 9b). More recently, differences in sea-level rise trends have stopped or  
476 even reversed. Between 1995-2017, rates were slightly larger in Jacksonville than Mayport ( $4.2$   
477  $\pm 1.5$  mm/yr. vs.  $3.7 \pm 1.3$  mm/yr), though results are not statistically different. Sea-level rise  
478 variations may in part reflect differences in subsidence; vertical land motion rates in the  
479 northeastern Florida region are -1 to -2 mm/yr, with considerable variability and uncertainty due  
480 to short GPS/GNSS record lengths (Blewitt et al., 2018).

481 The differences in the sea-level rise in Jacksonville and the estuary mouth region during the 20<sup>th</sup>  
482 century may also in part be driven by a reduction in the average surface slope in water level  
483 caused by channel deepening (Equation 8; Figure 10). We isolate the effect of river discharge by  
484 first removing oceanic variability, by subtracting the monthly averaged water level measured in  
485 Mayport from all gauge series. Next, the super-elevation caused by river discharge is obtained by  
486 comparing mean water levels during average discharge ( $200 \text{ m}^3/\text{s}$ ) with those during periods of  
487 no net discharge. A measureable vertical offset was found between Mayport and upstream  
488 stations under conditions of zero discharge (see supplemental information); for consistency, this  
489 offset was removed from all data in Figure 10, including those for which no discharge  
490 measurements are available.

491 The measured and modeled rise in mean water level caused by average river discharge is small,  
492 particularly under modern conditions (Figure 10). Semi-analytical and numerical model results  
493 agree reasonably well with each other, and suggest an approximate halving of the water level rise  
494 caused by average river flow (order 0.05-0.07 m decrease in water level in Jacksonville). Modern  
495 data is consistent with the semi-analytical model, and shows an approximately linear rise  
496 between the ocean and Jacksonville, to within data accuracy; the numerical model results show a  
497 slightly larger rise. Historical measurements from 1929-1932 are consistent with the historical  
498 model result, while the more uncertain data from 1892 is biased high. Most of the modeled  
499 change in fluvial effects between historical (blue) and modern (red) curves occurred near the  
500 ocean, seaward of Rkm 25-30; this is also where the largest increases in depth occurred (Figure  
501 3). Upstream of Rkm 25-30, the increase in average depth is less, likely leading to a smaller  
502 change in river slope (Equation 8). The doubling of tidal discharge (Figure 6b) through this  
503 section also tends to counteract the effect of depth increases (Equation 8).

504 Both measurement and modeling limitations likely influence results in Figure 10. For example,  
505 both the numerical and analytical model neglect sources of mean discharge below Jacksonville.  
506 Also, the analytical model does not consider the Stokes drift compensation flow caused by the  
507 correlation of horizontal and vertical tidal motions (see Moftakhari et al. 2016 for a definition;  
508 this discharge is estimated to be  $\sim 25\%$  of the mean flow in Jacksonville, based off of  
509 measurements). Other assumptions—such as the assumption of zero bed slope—could make a  
510 slight difference in the analytical model results. Further, neither the numerical or analytical  
511 model include wind effects or the mean slope caused by salinity intrusion. Many sources of  
512 precision and bias error add uncertainty to the measurements as well. The large variability  
513 around the mean, shown by the grey shading in Figure 10, shows that many processes—from  
514 wind to discharge—drive month-to-month variability. In Jacksonville, average water levels less  
515 than 1km from each other differ by 0.01-0.02 m (Figure 10). The reasons are unclear, but could  
516 include leveling error, benchmark or datum drift, differences in subsidence, or real differences in  
517 water level, for example, transverse water surface slope. Nonetheless, both measurements and



518 models support the inference that channel deepening has reduced the response of mean water  
519 level to increases in discharge.

520

### 521 3.3 Measurements during Hurricane Irma

522

523 During hurricane Irma (September 10-12, 2017), the maximum total water level (TWL) at both  
524 Mayport and Jacksonville (NOAA gauge) reached 1.7 m relative to the NAVD-88 datum.

525 However, the timing of the peak and the hydrodynamic factors contributing to the water level  
526 were different. At the estuary inlet, measured water level peaked slightly more than 2 hours after  
527 the predicted high tide of 0.64 m (Figure 11). By contrast, peak water levels in Jacksonville  
528 occurred on the following high tide. At Jacksonville, waters stayed within 0.05 m of peak TWL  
529 for 2.5 hours (Figure 11a), with the long duration likely contributing to the severity of flooding.  
530 At Mayport, water levels only briefly attained a peak and remained above 1.5 m for less than an  
531 hour.

532 We estimate that the predicted tide at Jacksonville was  $\sim 0.13$  m larger today than it would have  
533 been under historical conditions, given the approximately 90% increase in tidal range there  
534 (Figure 7). Hence, tides likely played a larger role in the total water level during Irma than they  
535 would have a century ago. Fortuitously, a worst case scenario—amplified tides occurring in  
536 phase with storm surge—was avoided. At the coast, storm surge (measured – predicted water  
537 level) peaked approximately half an hour before the predicted low tide (Figure 11). A similar  
538 timing occurred in Jacksonville. Hence, as the tide was rising in Jacksonville, the storm surge  
539 was falling, counteracting each other (Figure 11).

540 The long time scale of flooding at Jacksonville occurred because of the combined effect of storm  
541 surge, local discharge, and the astronomical tide. The local discharge wave peaked  
542 approximately 1 day after peak flood waters (Figure 11d); hence, the rising arm of the discharge  
543 hydrograph added significantly to the observed peak water level. The large precipitation of  
544 between 0.18-0.28 m of rain within Jacksonville (Cangialosi et al., 2018) likely influenced the  
545 local discharge. Another factor was the southerly (south-to-north) wind in the eastern quadrant of  
546 hurricane Irma as it moved north thru the western portion of Florida (domain B in Figure 1; see  
547 also Bacopolous et al., 2009, 2017 for investigation of local wind effects). This likely produced  
548 significant local wind setup near Jacksonville at or near the same time that marine-sourced surge  
549 was peaking. As the rising arm of the discharge freshet meets the storm surge and tide, a surface  
550 water level difference develops between the coast and Jacksonville (Figure 11a). Due to the  
551 larger channel depths and lower frictional resistance today, a smaller water-level slope may have  
552 been required to drain this water today, than historically. We next investigate this idea by  
553 evaluating simulations of hurricane Irma.

554

### 555 3.4 Simulations of Hurricane Irma: Historical vs. Modern

556



557 Simulations show that both maximum tidal amplitudes and maximum storm surge during  
558 hurricane Irma increased everywhere due to channel reconfiguration and deepening, relative to  
559 what they would have been in 1898 (Figure 12a and 12b). The increase in both tidal amplitude  
560 and storm surge amplitude is spatially variable, rising from zero (no change) at the coast to a  
561 maximum increase at Rkm 23-25 near Dames Point of  $\sim 0.16$  m and  $\sim 0.57$  m, respectively  
562 (Figure 12a and 12b). Similar to tides (Figures 6,7, and 12a), the difference between modern and  
563 historical surge diminishes further upstream (Figure 12b). As discussed above (section 3.3),  
564 these peak tidal and surge amplitudes were not phased together, diminishing their combined  
565 effect.

566 In contrast to tides and surge, the simulated super-elevation in water level caused by peak flood  
567 discharge (about 1 day after peak water level) decreases significantly between the historical (blue  
568 line) and modern configurations (red line, Figure 12c). The difference between historical and  
569 modern water levels expands from zero near the inlet to  $\sim 0.6$  m near Dames Point (Rkm 23-25),  
570 and remains fairly constant upstream to Jacksonville (Figure 12c).

571 At its peak, total water level (tide + surge + river flow) was simulated to be up to 0.2 m larger in  
572 the historical configuration, except at the estuary inlet (Figure 13). There, peak water levels  
573 were driven primarily by storm surge, and occurred  $\sim 10$  hours earlier than upstream (Figure 13a;  
574 see also Figure 11). Individually, the contributions of tide, surge, and river discharge (Figure  
575 13b,c, and d) to the peak total water level (Figure 13a) are similar to Figure 12, just of smaller  
576 magnitude. Because the maximum river flow, surge and tidal amplitudes occurred at different  
577 times, the worst-case scenario was avoided (compare Figure 12 and 13). Overall, the modeled  
578 peak water level agrees well with measurements (red-dots in Figure 13a), except for the  
579 anomalously low USGS measurement in Jacksonville.

580 Overall, the increase in marine-sourced water levels (tides+ surge) in the modern model is  
581 counteracted by a decrease in fluvial (river discharge) water levels (Figure 13, 14). Changes in  
582 both factors are small near the estuary inlet, but increase rapidly inland. The modeled increase in  
583 tides + surge is maximal in mid-estuary, and diminishes further upstream (Figure 14). By  
584 contrast, fluvial differences persist. Hence, the largest modeled decrease in peak water level  
585 ( $\sim 0.2$  m) from the historical configuration was simulated in Jacksonville; effectively, the sum of  
586 tide + surge effects ( $+0.25$  m) is less than river discharge effects ( $-0.45$  m; Figure 14). Based on  
587 Figure 13, approximately 10% of the decreased total water level in Jacksonville is attributable to  
588 the  $\sim 4$  km shortening in channel length to the ocean. The remainder is attributable to changes in  
589 depth, width, and drag coefficient. Overall, storm surge and tides contribute  $\sim 2/3$  to modern  
590 peak water levels, compared to  $\sim$ half under historical bathymetry.

591 Model sensitivity tests show that results remain qualitatively similar when the river discharge  
592 condition or the drag coefficient are altered (Figure 15). In these simulations, we leave oceanic  
593 forcing unchanged, but replace the virtual discharge condition with a constant river discharge  
594 (see section 2.5). Results suggest that for any fluvial discharge greater than  $\sim 2,600$  m<sup>3</sup>/s, the  
595 maximum water level in Jacksonville would have been higher, historically, than today (Figure  
596 15). For lesser discharge, the situation is reversed due to the effect of increased tides and storm  
597 surge. A simulated constant discharge of 3,500 m<sup>3</sup>/s and 6,000 m<sup>3</sup>/s produces modern water

598 levels that are consistent with USGS (1.54 m) and NOAA (1.69 m) peak measurements,  
599 respectively. Within this discharge range, simulated water levels in the historical configuration  
600 exceed modern levels, as in Figure 13. Because results in Figure 13 and Figure 15 are consistent,  
601 we surmise that the uncertainty involved in applying the modern discharge measurement as a  
602 virtual boundary condition in the historical configuration does not shift overall conclusions.  
603 Similarly, changing the Manning coefficient in the historical simulation modifies, but does not  
604 change, conclusions (see Figure 15).

605

## 606 4. Discussion

607

608 We next explore factors that help explain the simulated changes to tides, storm surge and  
609 extreme discharge, using the analytical models developed for typical (non-event) conditions.

610

### 611 4.1 Interpreting tidal and surge changes

612

613 The reasons for the spatially variable changes in tidal amplitudes (e.g., Figure 7) are next  
614 explored by simplifying the analytical solution to only consider the incoming wave (i.e., the  
615 amplitude  $A$  in Equation 4 is set to zero); this simplification can be made because the reflected  
616 wave is a small, order 10-20% correction except at bathymetric transitions. Further, we consider  
617 a constant width and depth section that is representative of near coastal bathymetry (Figures 2  
618 and 3), and artificially extend it upstream such that tides damp out. For explanatory simplicity,  
619 we assume that the linearized friction coefficient  $r$  is constant everywhere (Equation 3). Then,  
620 under highly frictional conditions in which  $\frac{r}{\omega} \gg 1$ , the tidal amplitude  $\eta(x)$  decays  
621 exponentially as

$$622 \eta(x) \approx \eta_o \exp(\mu x), \quad (10)$$

623 where  $\eta_o$  is the amplitude at the ocean boundary and the damping modulus  $\mu < 0$  is  
624 approximated as:

$$625 \mu \approx \frac{-2}{3} \left( \frac{\omega r}{gh} \right)^{\frac{1}{2}}. \quad (11)$$

626 Note that an additional dependence on depth  $h$  also enters through  $r$  (Equation 3); this solution is  
627 similar to LeBlond (1978). From Figure 6, the exponential decay in tidal amplitude is empirically  
628 estimated to be  $\mu = \frac{-1}{22} \text{ km}^{-1}$  (historical system) and  $\mu = \frac{-1}{38} \text{ km}^{-1}$  (modern system). The  
629 equivalent result is found in Equations 10-11 by reducing from  $\frac{r}{\omega} \sim 11.5$  (historical) to  $\frac{r}{\omega} \sim 7.7$   
630 (modern), using 5 m and 10 m as approximations for the depth of the lower 40km.

631

632 Following the observation that depth changes are the major cause of tidal amplitude changes  
633 (Figure 8), we next investigate how long-wave amplitudes in the simplified formulation

634 (Equation 10) depend on depth. Specifically, the change in amplitude  $\Delta\eta_H$  that occurs at point  $x$   
635 due to a change in depth  $\Delta h$  is approximated by taking the partial derivative of Equation 10 with  
636 respect to  $h$ , after substituting  $r = \frac{8}{3\pi} \frac{c_d U_T}{h}$  (Equation 3) into Equation 11. For simplicity we do  
637 not consider the changes in velocity  $U_T = \frac{Q_T}{bh}$  that occur due to deepening, and hold the velocity  
638 in Equation 3 constant. The modeled increase in tidal velocity of 20-35% in the lower 25km of  
639 the estuary is relatively small compared to the doubling of depth; moreover, the damping  
640 coefficient  $\mu$  (Equation 11) is more sensitive to depth variations ( $h$  dependence) than tidal  
641 velocity ( $U_T^{1/2}$  dependence). Hence, while increased tidal velocity is an important feedback effect  
642 that reduces the effect of depth changes, holding it constant is justified for scaling/interpretation  
643 purposes. Similarly, we neglect any small changes to tidal amplitudes at the ocean boundary  
644 caused by radiation damping, following the observation that M<sub>2</sub> amplitude changes in Mayport  
645 are slight ( $\sim 6\%$ ) compared to further upstream. We leave the effect of these and other factors  
646 (e.g., length and width changes) to future investigation.

647 With these simplifications, the change in tidal amplitude  $\Delta\eta_H(x)$  due to a depth change scales as:

$$648 \Delta\eta_H(x) \sim -\mu\eta_0 x * \exp(\mu x) \frac{\Delta h}{h}, \quad (12)$$

649 where the leading coefficient of order (1) has been dropped.

650 Despite the many simplifications and restrictions discussed above, several insights into the  
651 observed pattern of tidal change within the estuary follow from this analysis:

- 652 • Amplitude changes  $\Delta\eta_H$  are related to the percentage change in depth; in the Saint Johns  
653 River Estuary, depth increases have dominated historically over other modifications,  
654 since  $\frac{\Delta h}{h} \approx 2$  is quite large. The analytical dependency of tidal evolution on  $\frac{\Delta h}{h}$  also  
655 suggests that tides become progressively less sensitive to the same incremental change  
656  $\Delta h$  in deep waters, as opposed to shallow waters.
- 657 • The function  $x * \exp(\mu x)$  increases as one moves landward, reaches a maximum, and  
658 thereafter asymptotes towards zero. Hence, for an estuary described by the simplified  
659 model above, changes are predicted to be small at the estuary mouth ( $x=0$ ) and far  
660 upstream ( $x \gg -\frac{1}{\mu}$ ). In between, there is a location with *maximum sensitivity to altered*  
661 *depth*. Both in-situ and modeled results follow this pattern (Figure 6 and 7), and its  
662 influence is also observed in peak water levels (Figure 14). Similar to the Saint Johns  
663 River, other studies have also found that the maximum increase in tidal amplitudes is  
664 found within estuarine regions marked by a strong damping of tidal amplitudes. This  
665 distinguishes a highly damped estuary from an estuary with a total reflection, since in  
666 that case the maximum change often occurs near the head of tides (Winterwerp et al.,  
667 2013; Talke & Jay, 2020).
- 668 • The location of *maximum change*,  $x_{max}$ , occurs around the e-folding scale for damping,  
669  $L_{damping} = \frac{-1}{\mu}$ . In the Saint Johns Estuary, the maximum change in tidal amplitudes—  
670 i.e.,  $x_{max}$  --is located between 20-25 km from the coast. This is approximately equal to  
671 the observed e-folding scale for damping in historical tide data.

672

673 The simplifications in Equations 10-12 means that they only qualitatively approximate real  
674 behavior. More complex approaches (e.g., Li et al., 2016) are required to assess the effects of  
675 cross-sectional variability, off-channel storage, and other system features. Still, the spatial  
676 change in tidal amplitudes suggested by Equation 12 is not dissimilar to those noted in Figure 6  
677 and 7. Interestingly, storm surge changes appear to follow a similar pattern, with a simulated  
678 maximum that is nearly co-located with the position of maximum tide change (Figure 12b,  
679 Figure 14); more research is needed. A similar location for maximum change in total water level  
680 magnitudes was modeled by USACE (2014) for a 50y and 100y storm event and an increase in  
681 depth from a 12.2 m (40 ft.) to a 14.3 m (47 ft.) channel. Changes to tidal amplitudes and the  
682 scaling in Equation 12 may therefore provide insights into the spatial pattern of storm surge  
683 changes (see also Familkhalili et al., 2020).

684

#### 685 4.2 Interpreting subtidal change

686 The analytical model for mean water level is consistent with the changes simulated by the  
687 Delft3D numerical model at low flow (Figure 10), and provides insights into the factors that may  
688 influence subtidal change. Specifically, the subtidal slope term,  $\alpha \sim \frac{C_d Q_R U_T}{g b h^2}$  (Equation 9),  
689 suggests that factors such as increased depth and decreased drag coefficient may reduce the  
690 modeled subtidal water levels. These factors appear to be partially counteracted by decreased  
691 width (through diking of wetlands) and increased tidal discharge and tidal velocity. Further, the  
692 observed shift in subtidal water levels is a function of discharge,  $Q_R$  (Figure 13). Therefore, any  
693 change in water levels caused by river discharge may become more prominent during river flood  
694 conditions (see Figure 12-15), though overbank flow effects must be considered (see e.g. Helaire  
695 et al., 2019).

696 The approximate agreement between Equation 9 and empirical measurements (Figure 10)  
697 highlights the role that tidal velocity may have in setting mean water levels. The semi-analytical  
698 tide model suggests that a 20-35% increase in tidal velocity occurred between the historical and  
699 modern configurations between 0-30km from the estuary mouth. In the upper Scheldt, historical  
700 trends in tidal velocity also served to increase the subtidal slope in water level (Wang et al.,  
701 2019). The role of tidal velocity contrasts with large river systems such as the Mississippi, in  
702 which tides are small and thus neglected in models of the subtidal water level curve (e.g.,  
703 Nittrouer et al., 2012).

704 A qualitatively similar decrease in mean water levels due to channel deepening has been  
705 observed or modeled in other systems, including the Hudson River (Ralston et al., 2019), the  
706 Columbia River (Jay et al., 2011; Helaire et al., 2019), and the Ems River (Jensen et al, 2003).  
707 These observations were made between 100-250 km from the open coast, where the integrated  
708 effect of small changes in the slope of surface water level becomes more obvious. The modeled  
709 drop in water level between Jacksonville and the estuary inlet is small (Figure 10), except during  
710 extreme floods (Figure 15), and is qualitatively consistent with available in-situ data during  
711 average discharge conditions (Figure 10).

712 The flooding caused by river discharge during hurricane Irma is a low probability event. Only  
713 one other river discharge event besides hurricane Irma exceeded 3000 m<sup>3</sup>/s since 1988 (Sept.  
714 2004). More than half of the annual peak discharges measured in Jacksonville range between

715 1,000 and 1,600 m<sup>3</sup>/s. Therefore, the large effect of bathymetric changes on the water levels  
716 induced by fluvial effects during Irma (Figure 12-15) is unusual. Under more typical discharge  
717 conditions, a storm surge with a comparable magnitude to hurricane Irma would likely produce a  
718 larger water level today, than historically, particularly if phased together with tides (Figure 15).

719

#### 720 4.3 Comparison with other studies and sources of uncertainty

721 Our results are generally consistent with past modeling efforts. The large effect of fluvial forcing  
722 is consistent with Bacopoulos et al., (2017), who found that run-off from Tropical Storm Fay  
723 (2008) added ~0.5 m to the simulated storm tide. Similarly, modeling has suggested that the  
724 stormtide (surge + tides) measured in/near the shipping channel during a 50y and 100y event  
725 would increase by up to 0.2 m, after deepening from 12.2 to 14.3 m (40 to 47 ft.; USACE, 2014).  
726 Nonetheless, since the timing of a storm surge relative to tides and the river hydrograph may  
727 shift in each event, the modeled response to system changes may vary. For example, Bilskie  
728 (2013) found a negligible change in total water levels when hurricane Dora was modeled for  
729 both 12.2 and 14.3 m channels. As shown by Familkhalili & Talke (2016), greater tide  
730 magnitudes can reduce or negate the effect of an amplified storm surge, if the storm peak is  
731 timed at low water. Therefore, an approach that considers a full range of different storm tracks,  
732 magnitudes, rainfall, and tides is likely needed, to fully assess changes to flood hazard caused by  
733 channel deepening (e.g., Orton et al., 2020).

734 Our approach yields reasonable results that explain changes to empirically measured tides.  
735 However, the interaction of estuary tides with the open ocean can produce changes at the ocean  
736 boundary (e.g., radiation effects) which we do not consider analytically. Similarly, a drawback of  
737 the storm-surge hydrograph method used here is that storm surge magnitudes can vary along the  
738 ocean boundary due to meteorological forcing (e.g., Dietsche et al., 2007). The (likely small)  
739 errors that are introduced by assuming a constant storm surge elevation along the ocean  
740 boundary are present in both the historical and modern models, and therefore have little effect on  
741 their comparison.

742 The good correspondence between analytical and numerical results (for tides and mean water  
743 levels) under average conditions suggests that channel deepening is the major cause of changed  
744 numerical results; however, this inference has not been rigorously tested with one-at-a-time  
745 variations in numerical model bathymetry and forcing (depth, wetland connectivity, surge  
746 variability, etc). Moreover, the analytical result represents a simplified system with idealized  
747 bathymetric variation and no wetlands; these factors may explain why the friction coefficients  
748 used in the analytical model were larger than in the numerical model. Many additional factors  
749 have not been considered. For example, we do not directly model changes to local setup caused  
750 by local wind, though these may to some extent modeled through using a virtual discharge  
751 boundary condition (see Section 2.5). Many additional factors beyond channel deepening and  
752 shortening likely influence flood heights in Jacksonville. For example, we do not analyze the  
753 effect of the entrance jetties. Over the past century, the watershed around Jacksonville has  
754 become more urban, and natural streams have been channelized, both tending to make the run-  
755 off response more immediate. We do not consider the effect of such land-use changes. Other  
756 factors, such as barometric pressure variations and salinity intrusion, also influence water level  
757 patterns (Bacopoulos et al., 2009, Orton et al., 2012, Mulamba et al., 2019). Nonetheless, the

758 simulated decrease in river discharge effects during hurricane Irma (Figures 12-15) is consistent  
759 with observations of decreased water level during low flow conditions (Figures 9-10) and with  
760 analytical scaling (Equation 9). Similarly, the amplification in simulated surge is consistent with  
761 observations and analytical models of tides. Hence, our results suggest a substantial change in  
762 barotropic dynamics within the Saint Johns River Estuary, with corresponding effects on flood  
763 hazard.

## 764 5. Conclusions

765 In this contribution, we investigate how channel deepening, shortening and other modifications  
766 alter the way tides, storm surge, and river discharge flow through hyposynchronous estuaries  
767 marked by tides that strongly decay in the landward direction. The results suggest that long-  
768 wave amplitudes in estuarine regions marked by strong damping are quite sensitive to changes in  
769 depth. These changes manifest in a spatially variable way, with a maximum that is located near  
770 the observed damping length-scale for tides. Subtidal water levels, by contrast, are predicted to  
771 decrease due to the same channel deepening. These predictions are tested in the Saint Johns  
772 River Estuary, Florida, an estuary in which depths have approximately doubled, width decreased,  
773 and the shipping channel shortened since the 1890s. Both in-situ, numerical, and analytical  
774 results indicate that tidal amplitudes and tidal discharges have increased, and in many locations  
775 doubled, in response to channel dredging and to a lesser extent width, length, and drag  
776 coefficient changes. Storm surge has also increased. Nonetheless, modeled subtidal water levels  
777 have decreased, particularly during extreme flood flows. As a result, hurricane Irma likely  
778 would have caused higher water levels, had it occurred in 1898.

779 Since many estuary regions are highly frictional and marked by a strong damping in tides (e.g.,  
780 Talke & Jay 2020 review, and references therein), the spatially variable changes to tides and  
781 surge observed in the Saint Johns River Estuary likely occurs in other locations. An implication  
782 is that flood hazard may be shifting in a spatially non-uniform way over time, due both to  
783 changes in long-waves and subtidal water levels. As was also observed by Ralston et al. (2019)  
784 for Albany, New York, larger tides and storm surge magnitudes in a modern system can  
785 paradoxically be correlated with less flooding than would have occurred historically, at least for  
786 the event considered here. Nonetheless, Jacksonville is probably more at risk to flooding from  
787 large hurricane surge than it was historically. Effectively, as in the Cape Fear Estuary (see  
788 Familkhalili & Talke, 2016), the natural protection afforded by shallow channels has been  
789 largely removed, making inland regions much more exposed to marine-sourced flooding. In  
790 estuaries and tidal rivers, therefore, studies that evaluate changing flood hazard must consider the  
791 (often competing) sum of river, tidal and surge effects.

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793

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797 NOAA, or the EV2 database (<https://www.ncdc.noaa.gov/EdadsV2/>). The other archival tide



798 records are available in Record Group 23 at the US National Archives in College Park, Maryland  
799 (accession number RG 23, UD-WW Entry 14). Selected pictures and descriptions of important  
800 archival data and meta-data are included in the Supplemental Information.

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- 970

971 Tables

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973

974 *Table 1: Comparison of measured and modelled M<sub>2</sub> tidal statistics at downtown Jacksonville.*  
 975 *Water level phase is defined relative to the value in Mayport (Rkm 5.5). Relative phase is*  
 976 *defined as  $\varphi_h - \varphi_Q$ , where  $\varphi_h$  and  $\varphi_Q$  are the phase angle of the vertical tide and cross-*  
 977 *sectionally averaged discharge, respectively. Historical estimates of discharge amplitude are*  
 978 *obtained from Rkm 30 (see Figure 6b). D3D = Delft 3D numerical model.*

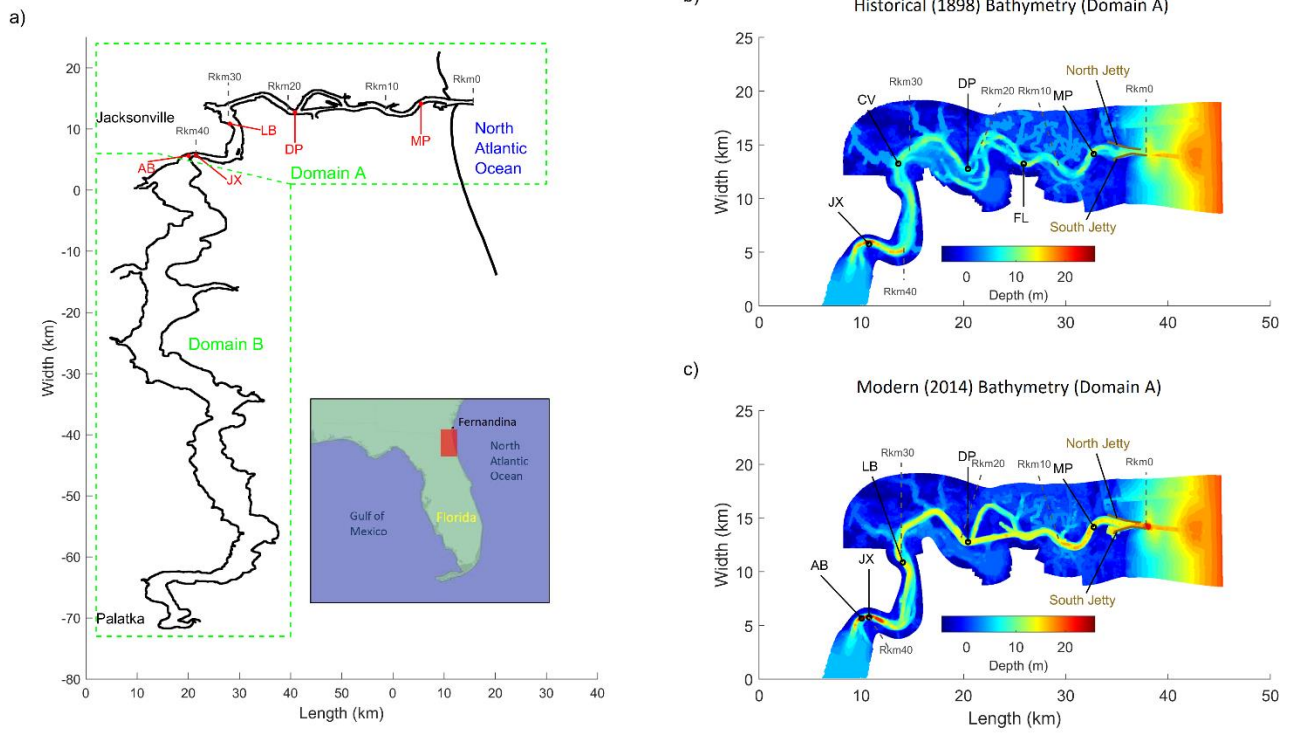
		RMSE for tidal amplitude (m)	Tide Phase Difference (Jacksonville – Mayport) (degrees)	Relative Phase between tidal discharge and amplitude (degrees)	Tidal discharge amplitude (m <sup>3</sup> /s)
Modern	Measured		49	3	4200
	D3D	0.008	41	-8.6	4130
	Analytical model	0.035	30	1.5	3800
Historical	Measured		40	Not known	2050
	D3D	0.025	64	5	2400
	Analytical model	0.044	44	-2	1850

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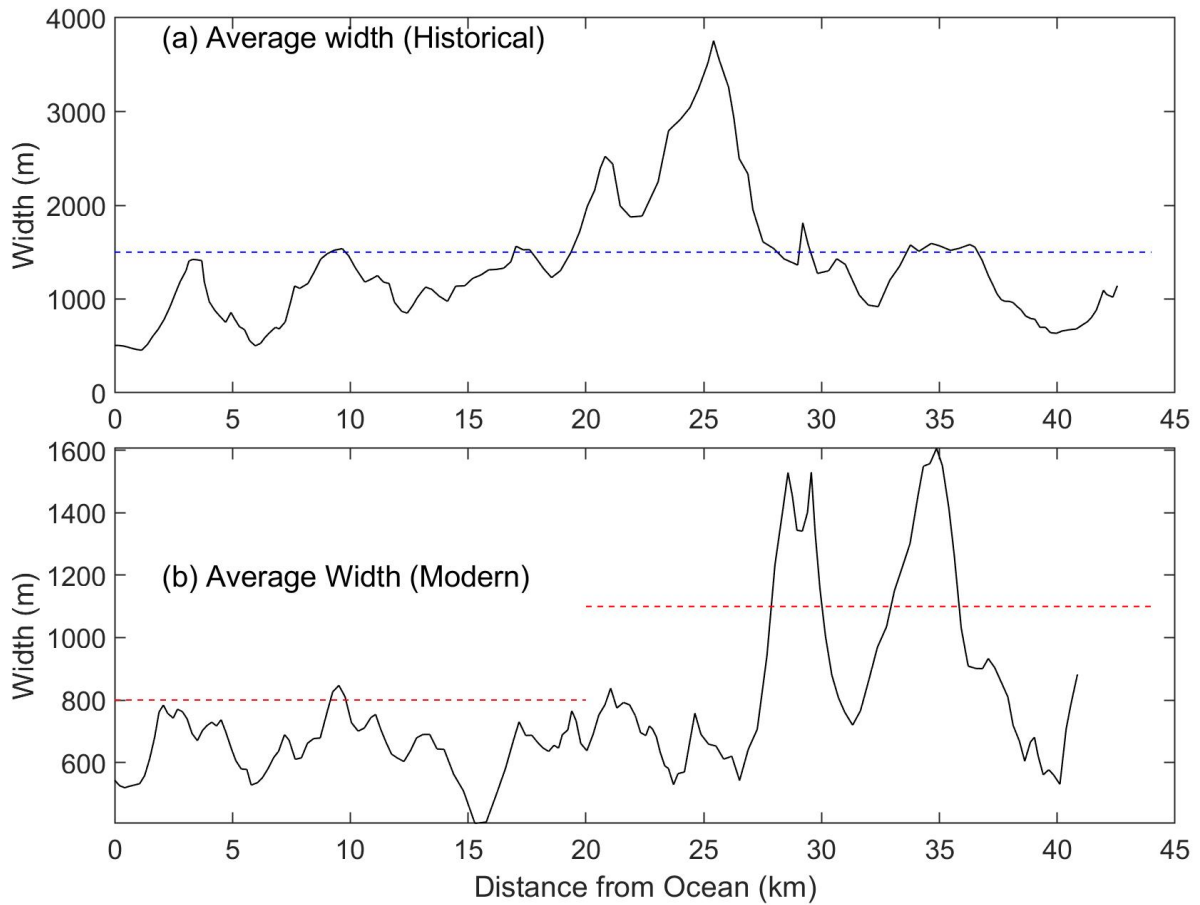


981 Figures



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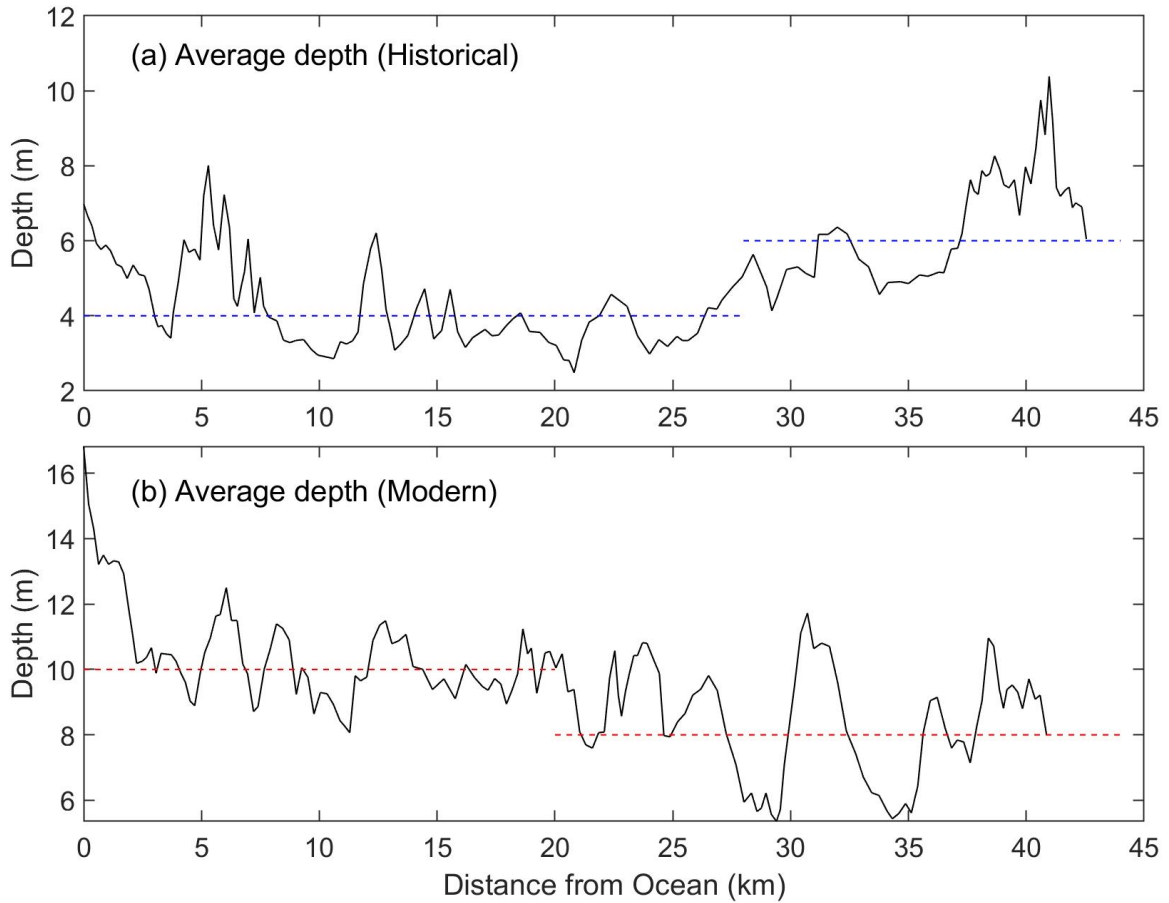
983 *Figure 1: a) Site map of the Saint Johns River Estuary, Florida, with b) Historic (1898) and c)*  
 984 *Modern (2014) bathymetry depicted from the ocean to Jacksonville. Abbreviations as follows:*  
 985 *AB = Acosta Bridge, USGS gage 02246500, JX = Jacksonville, NOAA gage 8720226,*  
 986 *LB=Long Branch, NOAA gage 8720242, DP =Dames Point, NOAA gage 8720219, MP =*  
 987 *Mayport, NOAA gage 8720218.*



988

989 *Figure 2: Measured width in (a) 1898 and (b) 2014 bathymetry, for a cross-section that extends*  
 990 *between the MLW datum located on either side of the channel thalweg. The dashed line indicates*  
 991 *the depth used in the idealized tide and river discharge models.*

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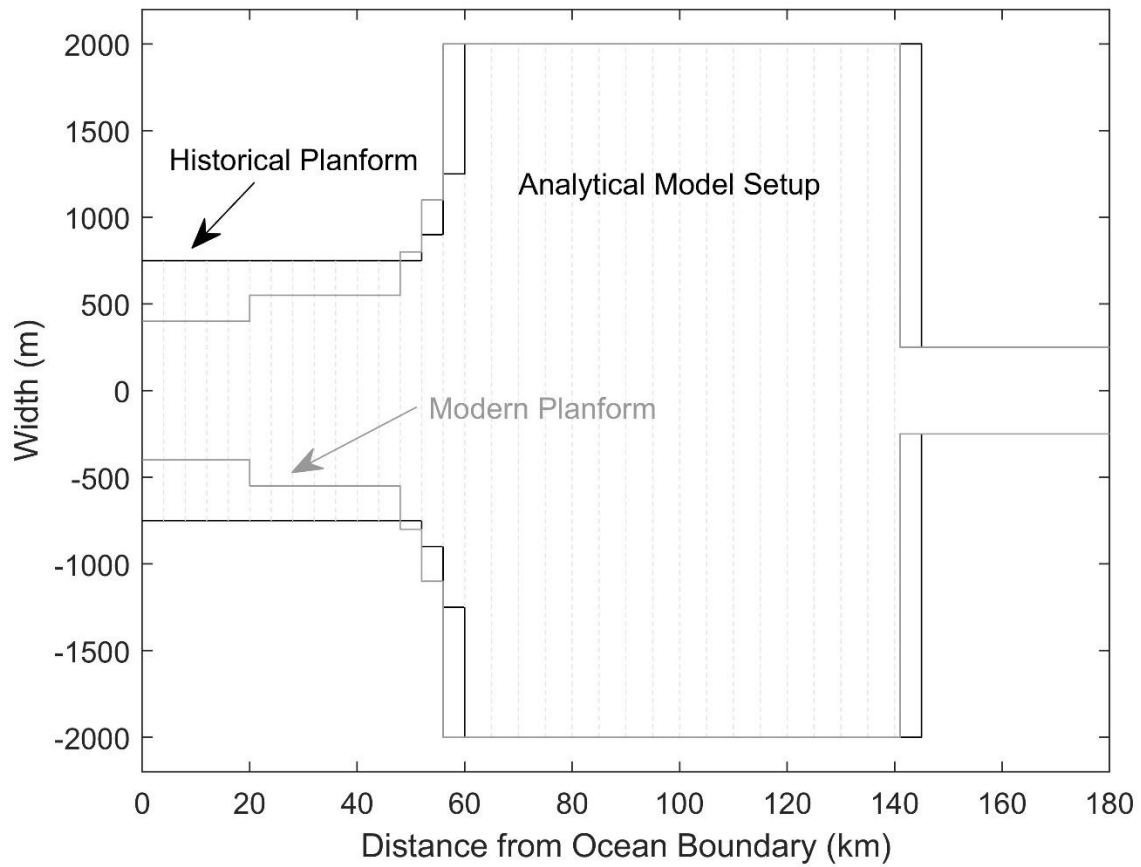
994 *Figure 3. Average estimated depth in (a) 1898 and (b) 2014 bathymetry, obtained by dividing*  
 995 *the cross-sectional area by the cross-sectional width. Datum is mean sea-level. The dashed line*  
 996 *indicates the depth used in the idealized tide and river discharge models.*

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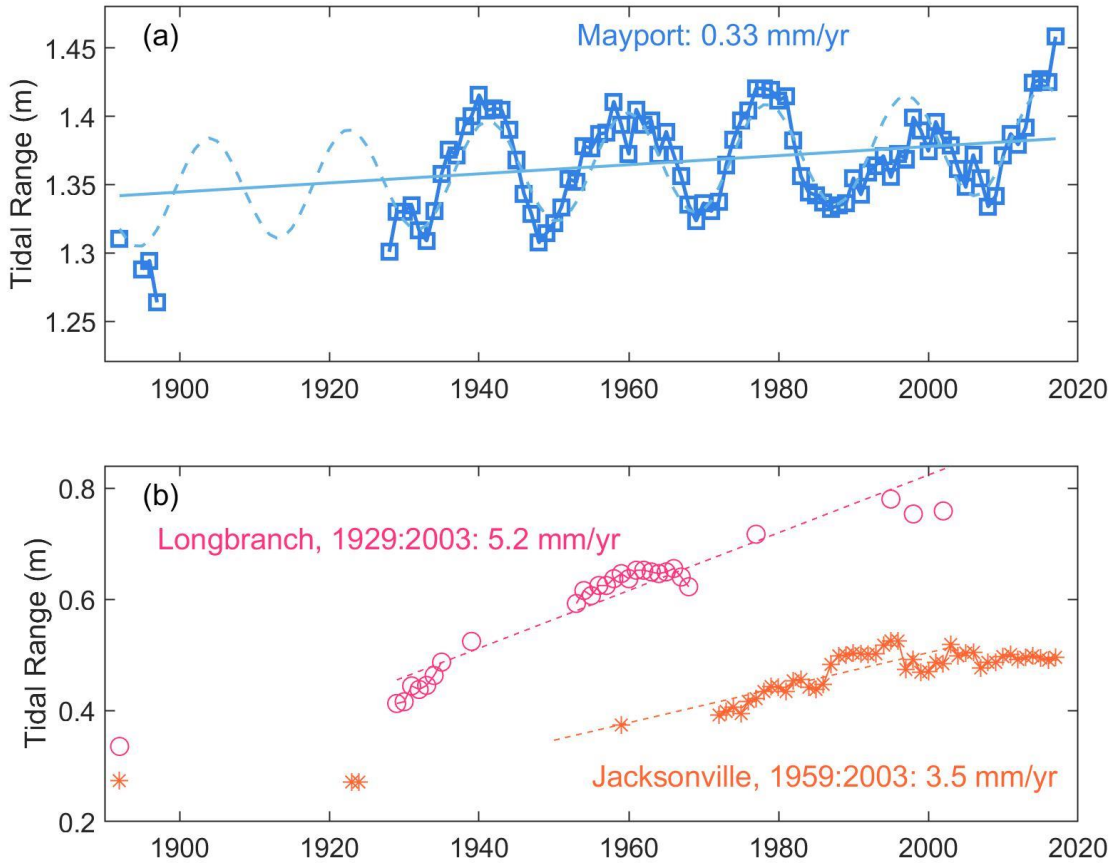
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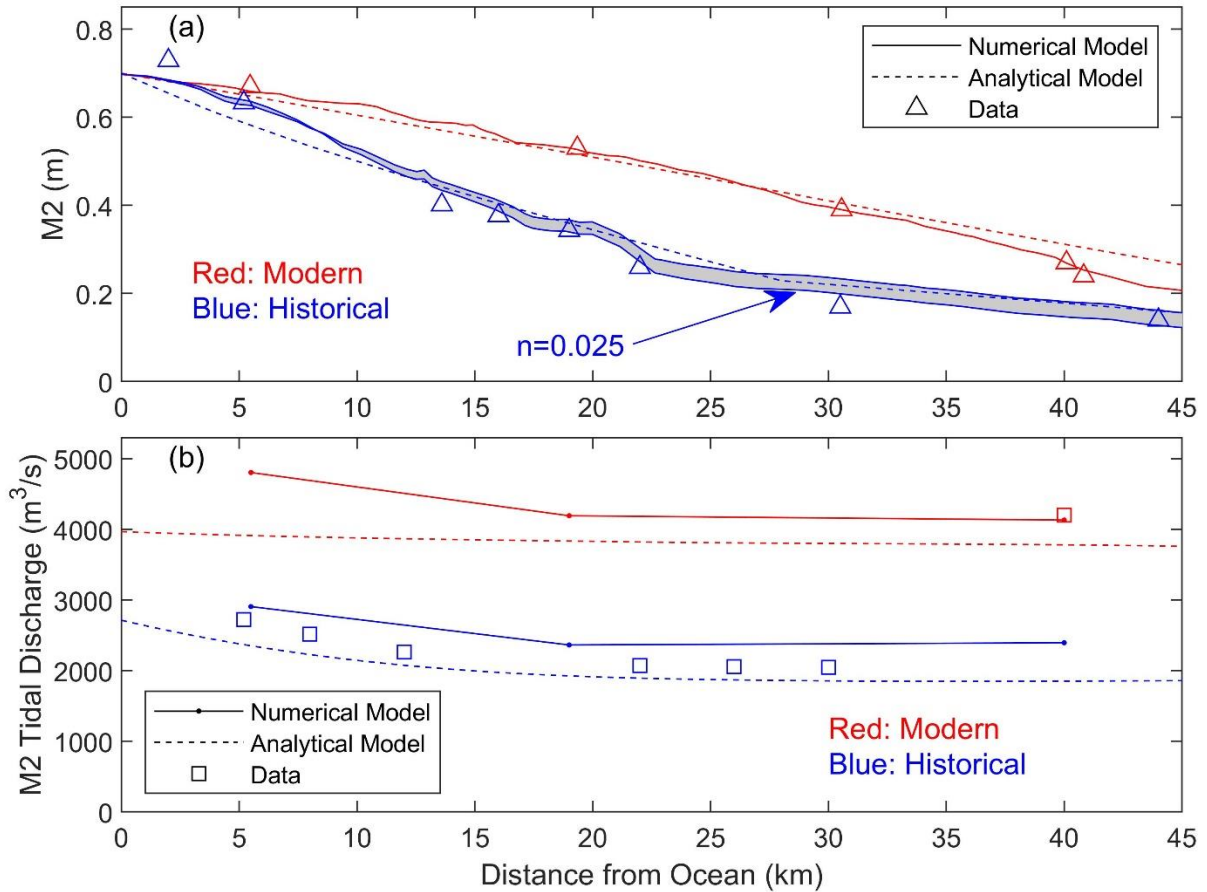
1002 *Figure 4: Planform of the idealized tidal channel model developed in section 2.3, for both the*  
 1003 *historical and modern configurations. The channel at the right hand side extends an additional*  
 1004 *~100km to enable the tide to damp out.*



1005

1006 *Figure 5: Changes in annual mean tidal range since the 1890s at (a) Mayport (~River km 5.5)*  
 1007 *and (b) Jacksonville-Longbranch (River km 31, violet color) and downtown Jacksonville (River*  
 1008 *km 40, orange color). Trends were obtained using robust linear regression and had a standard*  
 1009 *error of 0.06 mm/yr. (Mayport), 0.3 mm/yr. (Jacksonville-Longbranch) and 0.4 mm/yr.*  
 1010 *(Jacksonville).*

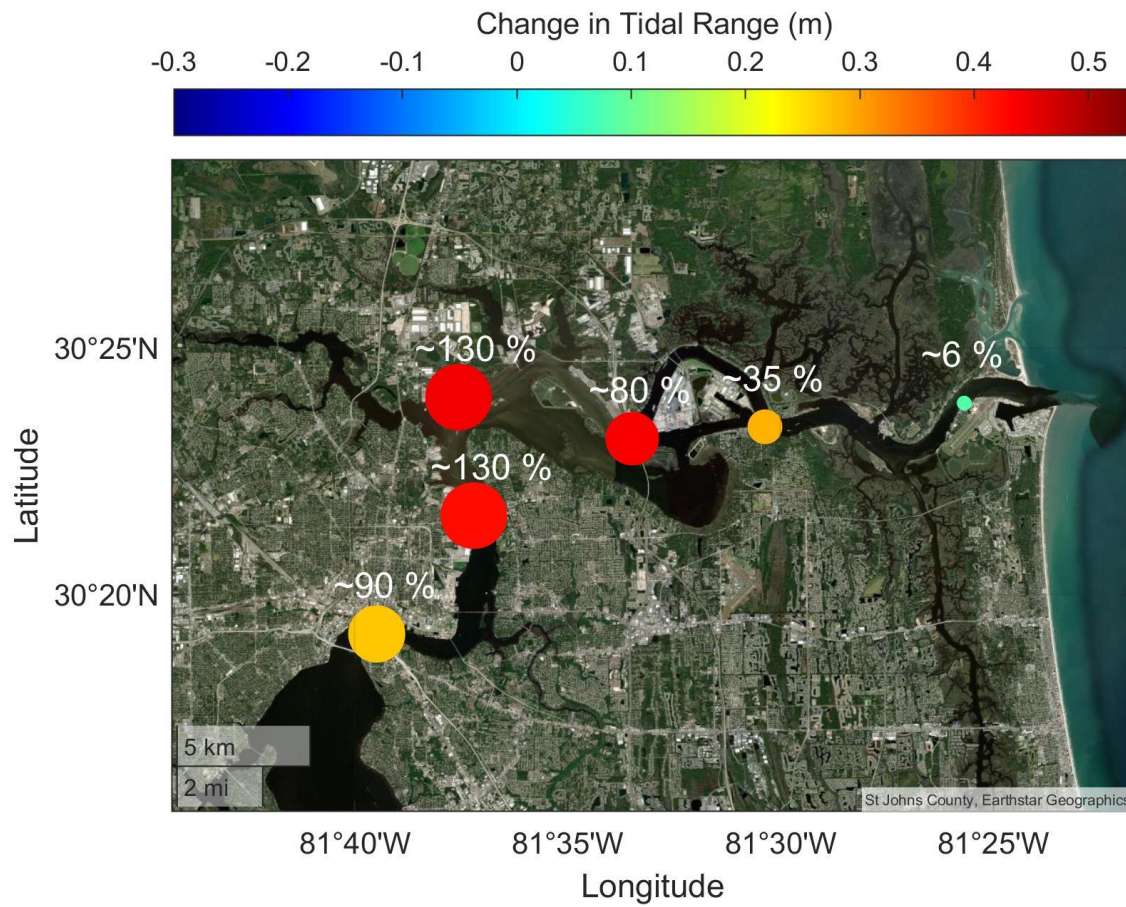
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1013 *Figure 6: Comparison of historical and modern  $M_2$  amplitude (a) and  $M_2$  tidal discharge (b). In*  
 1014 *(a), the grey shading bounds the numerically modelled tidal amplitudes that occur for a*  
 1015 *Manning's  $n$  value of  $0.025 \text{ s/m}^{1/3}$  (bottom line) vs.  $n = 0.02 \text{ s/m}^{1/3}$  (top line). Historical tidal*  
 1016 *and discharge estimates primarily from Gieseler (1893), with a few additional tidal amplitudes*  
 1017 *obtained from archival Coast and Geodetic Survey records (see Supplemental Information).*

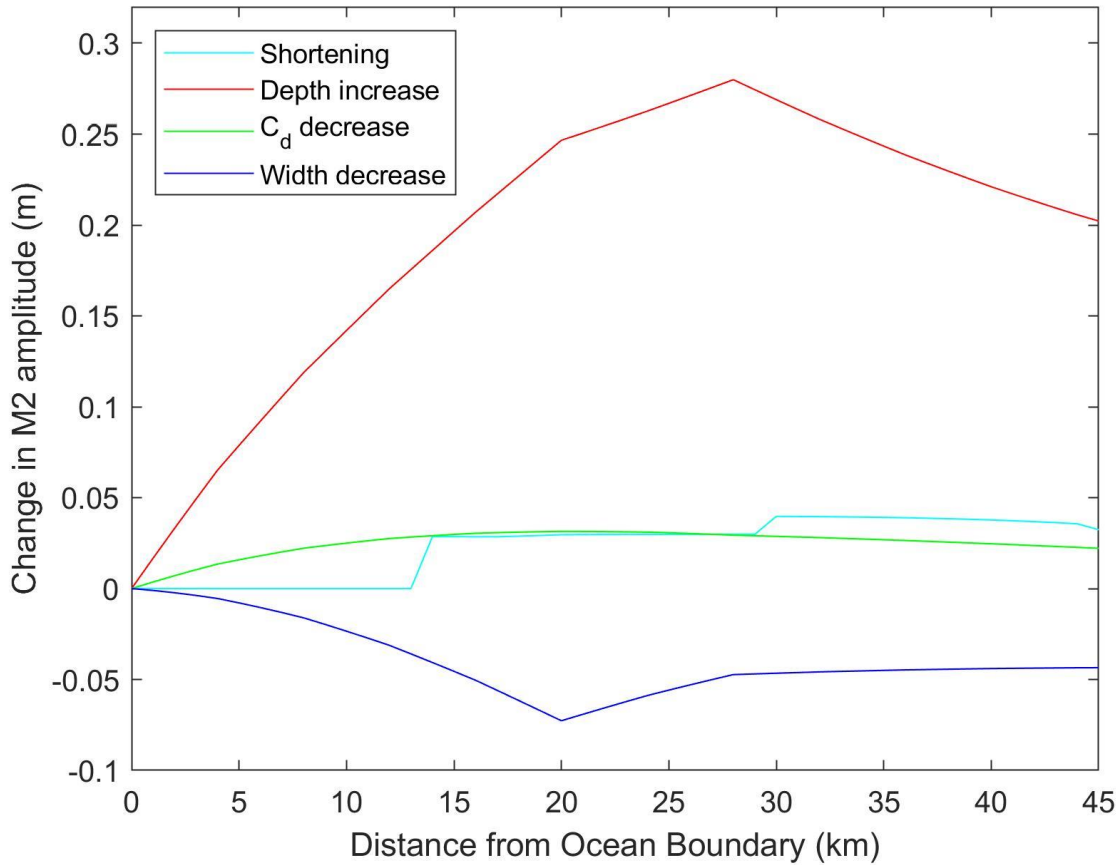
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1020 *Figure 7: The spatial change in tidal range, based on modern minus historical values (see*  
 1021 *Figure 6). The size of each bubble is proportional to the total change. The percent increase*  
 1022 *relative to historical conditions is indicated.*

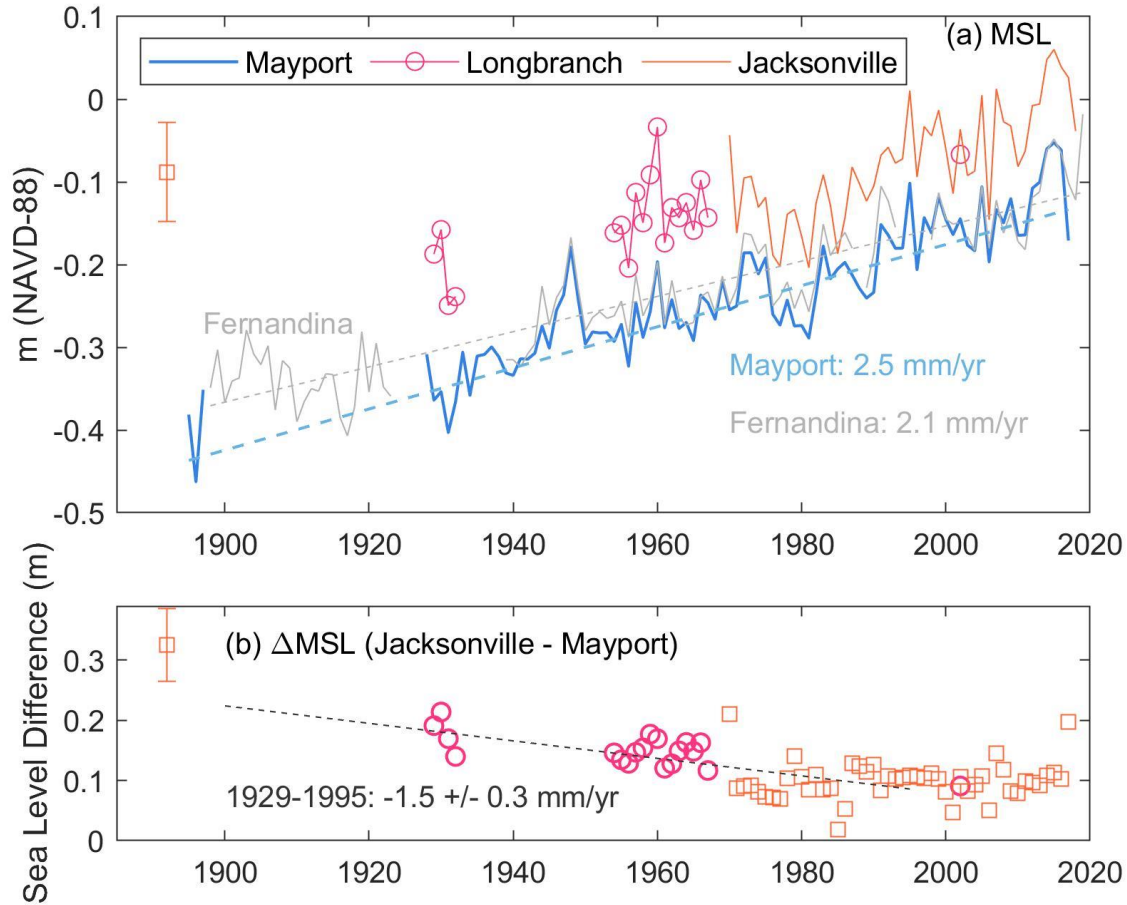




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1024 *Figure 8: Causes of tidal change, estimated using the analytical tide model. Results obtained*  
 1025 *using the modern configuration, changing one parameter at a time. Results show how much*  
 1026 *tidal amplitude would change if one component was changed from the historical value to the*  
 1027 *modern value. For example, increasing depths from the historical to the modern value (see*  
 1028 *Figure 3) would produce a maximum increase of nearly 0.28 m in tidal amplitude.*





1029

1030 *Figure 9: (a) Comparison of annually averaged mean sea level at Mayport (coastal station) and*  
 1031 *at Longbranch (eastern Jacksonville, Rkm 30) and downtown Jacksonville (~ Rkm 40) (b) Since*  
 1032 *the 1890s, the annual difference in mean sea level between Jacksonville and Mayport ( $\Delta$ MSL)*  
 1033 *decreased. In (a), a difference in sea-level rise of ~0.4 mm/yr. is found between the gauges at*  
 1034 *Mayport and Fernandina (grey line), located 30 km north.*

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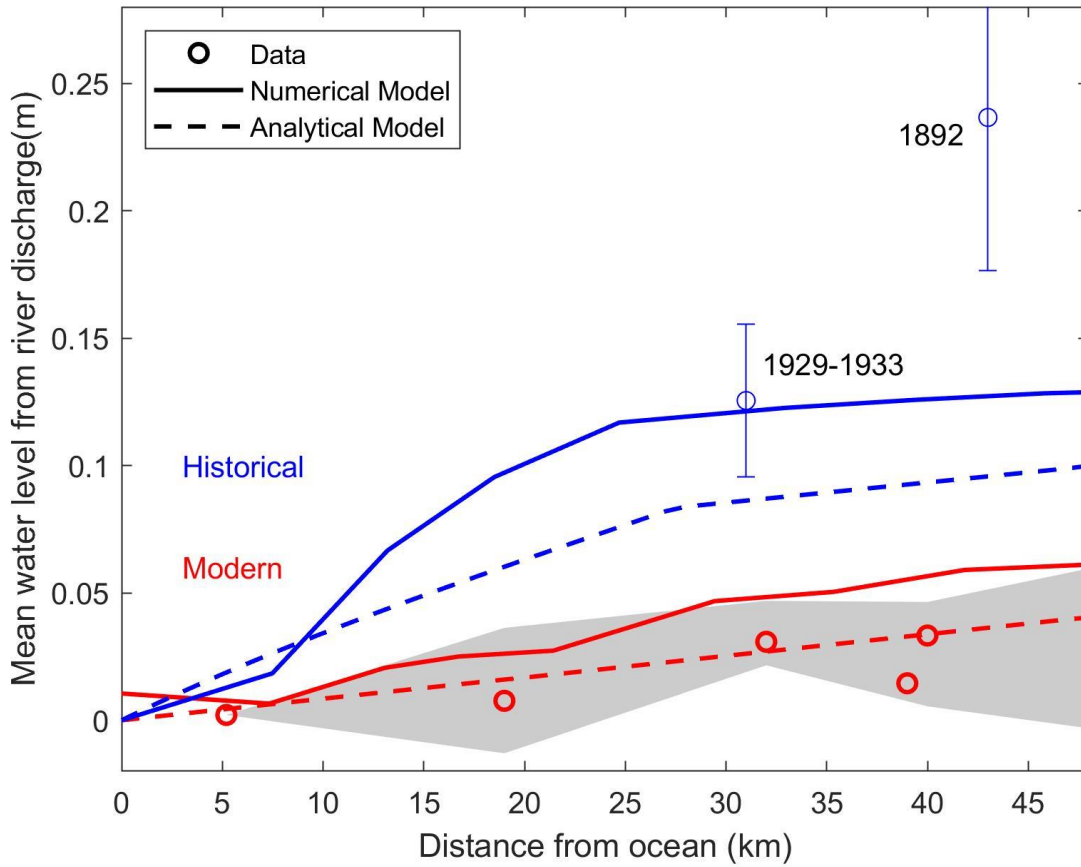
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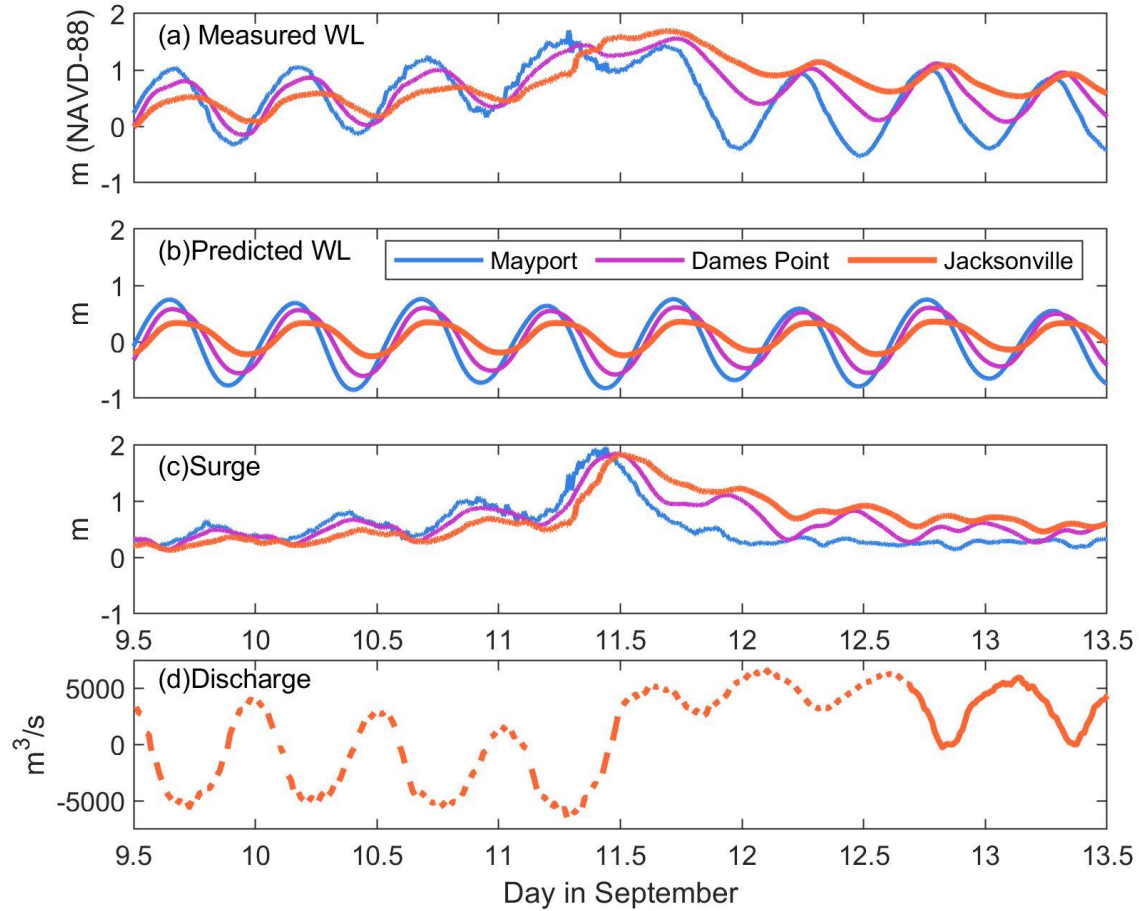
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1046 *Figure 10: Comparison of modeled and measured mean water level caused by river discharge,*  
 1047 *for both historical (blue) and modern (red) conditions. Data are based on the difference*  
 1048 *between monthly averaged water level between a station and the monthly water level in Mayport*  
 1049 *(Rkm 5.5). For each location, some residual difference in water level occurred at zero discharge*  
 1050 *in modern measurements; this offset was removed from both modern and historical*  
 1051 *measurements. The shaded region depicts the 10<sup>th</sup> and 90<sup>th</sup> percentile of measurements.*

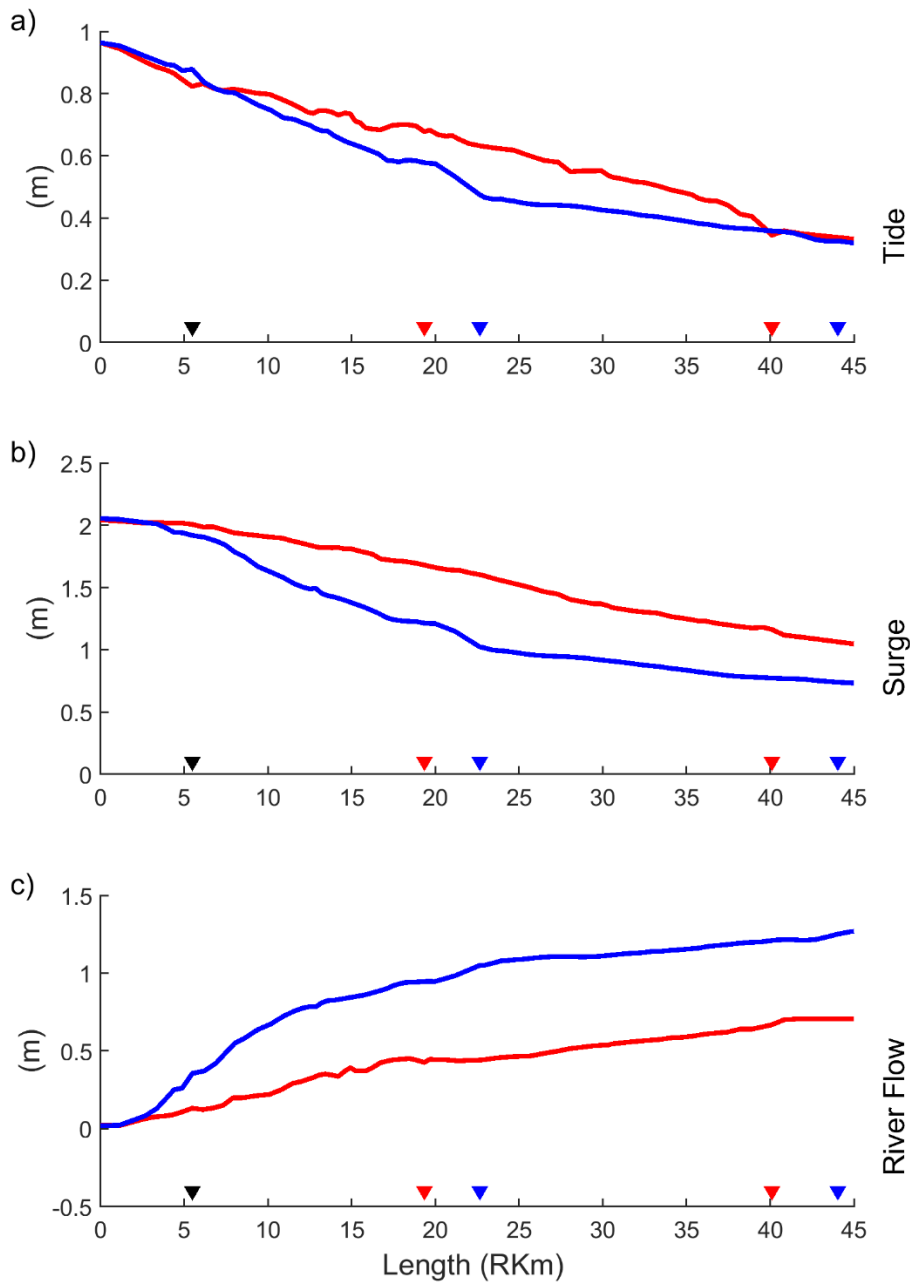
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1054 *Figure 11: Water Level and Discharge in the Saint Johns River Estuary in September 2017. (a)*  
1055 *Measured Water Levels; (b) Predicted tidal water levels (from NOAA); (c) The difference*  
1056 *between measured and predicted water levels, or surge; (d) the measured discharge in*  
1057 *Jacksonville (blank spaces denote recording gaps). Data are from Mayport (Rkm 5.5; blue),*  
1058 *Dames Point (~Rkm 19; purple) and Jacksonville (Rkm 40; orange); see legend in (b).*

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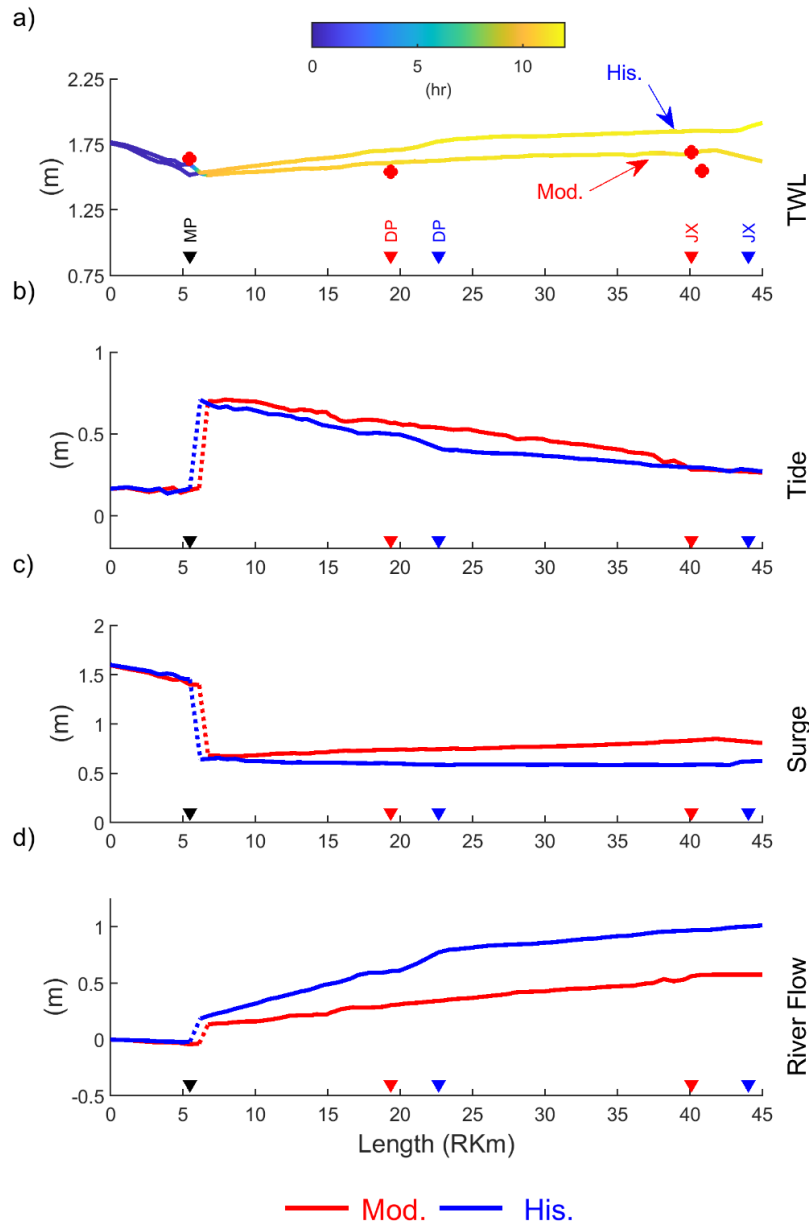


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1061 *Figure 12: Simulated peak amplitude of (a) tides; (b) storm surge and (c) river flow effects*  
 1062 *during hurricane Irma. These peaks occurred at different times and did not coincide with the*  
 1063 *overall maximum water level. Coloring denotes the modern configuration (red) and historical*  
 1064 *configuration (blue). Datum is the still water level in the ocean domain.*

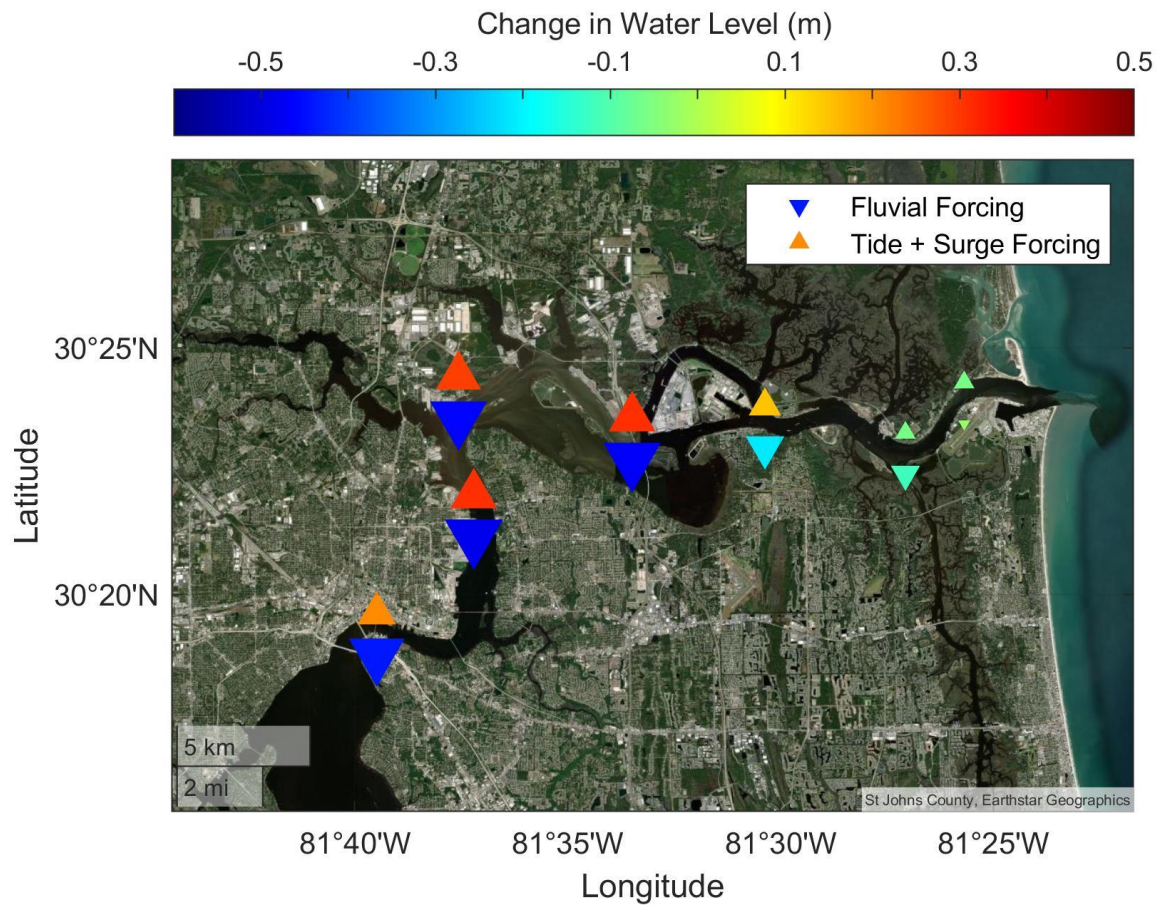
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1067 *Figure 13: (a) Maximum total water level simulated during Hurricane Irma for the historical*  
 1068 *(blue) and modern (red) configuration, with the contribution of tides (b), storm surge (c) and*  
 1069 *river flow (d) to the total water level shown below. Adding (b), (c), and (d) yields the total water*  
 1070 *level (a). In (a), the color shading shows how much the time of peak water level lagged the time*  
 1071 *of peak water level at the coast. The stair-case pattern in (b), (c), and (d) results from shifts in*  
 1072 *the timing of the peak. After a time shift, the relative contributions of tidal, surge and river*  
 1073 *forcing to the peak water level have changed. In (a) MP= Mayport (black triangle); DP =*  
 1074 *Dames Point; and JX = Jacksonville. The red and blue triangles mark the modern and historical*  
 1075 *distance of Mayport and Jacksonville from the coast.*

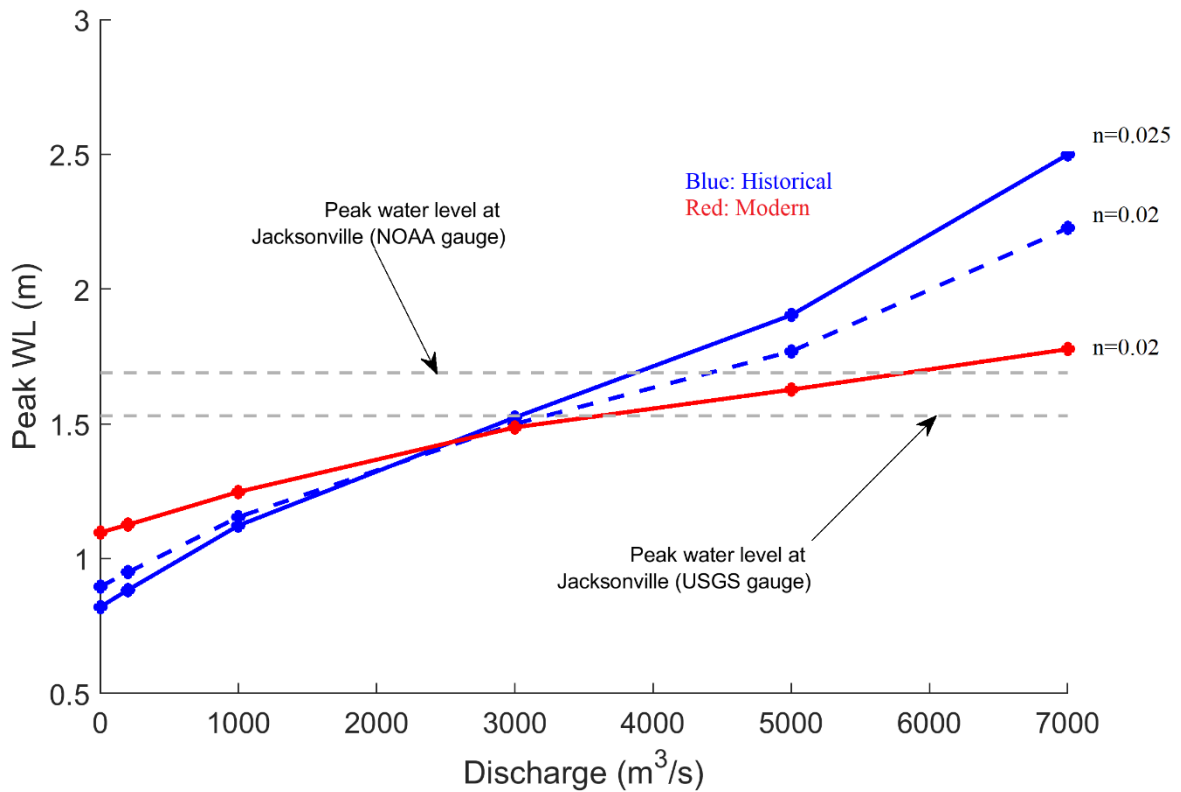
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1078 *Figure 14: Simulated change to peak water level (historical- modern) caused by fluvial forcing*  
 1079 *(river discharge) and marine (tide+ surge) forcing, based on Figure 13. The size of each symbol*  
 1080 *is proportional to the magnitude of the effect. Changes to fluvial and marine induced-water*  
 1081 *levels counteract each other.*

1082



1083

1084 *Figure 15: Simulated maximum water level at Jacksonville (Rkm 40) during hurricane Irma*  
 1085 *using constant discharge at the upstream boundary, for both 1898 (Historical) and 2014*  
 1086 *(Modern) bathymetry. The coastal boundary condition is the same as in Figure 12-14. The*  
 1087 *measured peak water level observed at the NOAA and USGS gauges in Jacksonville (1.69 m and*  
 1088 *1.54 m relative to the NAVD-88 datum, respectively) is indicated. The impact of changing the*  
 1089 *Manning's coefficient in the historical simulation is depicted by a solid line (n =0.025) and a*  
 1090 *dashed line (n = 0.02). Above a discharge of ~ 2,600 m<sup>3</sup>/s, water levels in Jacksonville in the*  
 1091 *historical simulation exceed the modern simulation.*